



Experimental aspects of nuclear (giant) resonances

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Normandie Université

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- 4 Example of experiments with unstable nuclei
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 - MSU
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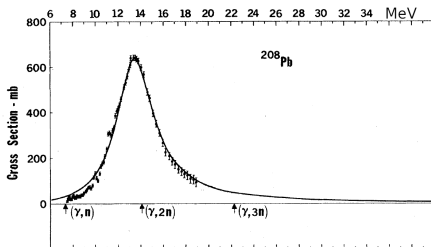
Giant Resonances Properties

Giant Resonances

Discovered in 1937 Bothe et al. Z. Phys. 71 (1937) in photo-absorption ;

Theoretically described in 1944 Migdal J. Phys, (USSR) 8 (1944)

Berman et al. Rev. Mod. Phys. 47 (1975)



Described/fitted by a Lorentzian
(Breit-Wigner)

$$\sigma_{\gamma}(E) = \frac{\sigma_{\max}}{1 + \left[\frac{E^2 - E_r^2}{E\Gamma_r} \right]^2}$$

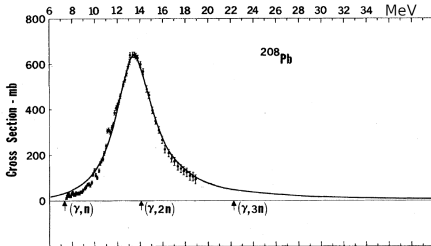
Complex plane: (E, Γ)

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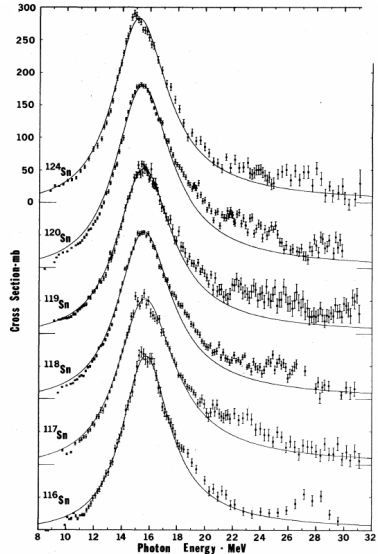
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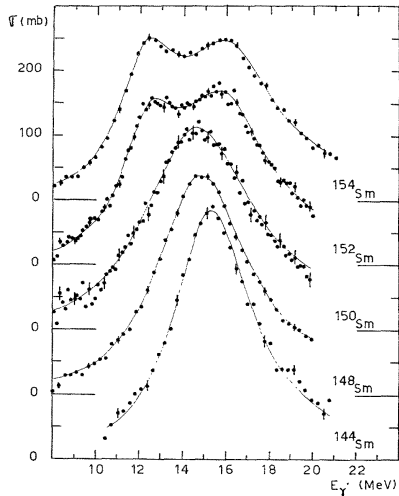
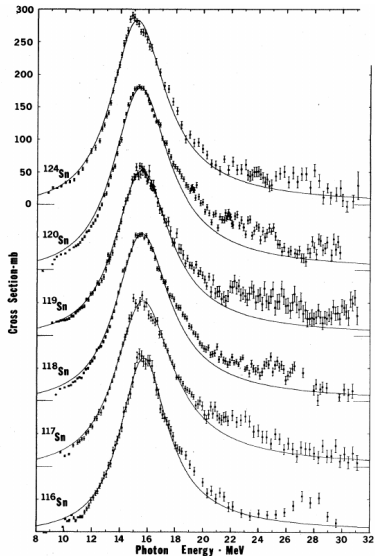
Complex plane: (E, Γ)



Giant Resonances

Discovered in 1937 [Bothe et al. Z. Phys. 71 \(1937\)](#) in photo-absorption ;

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Deformation [Carlos et al. Nuclear Physics A 225 \(1974\)](#)

Giant Resonances

Hydrodynamic models

Giant Resonances

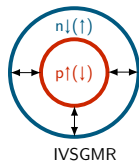
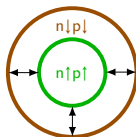
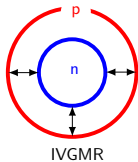
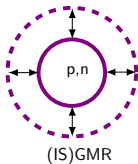
Giant resonances are high-frequency collective excitation of atomic nuclei

Macroscopic/Hydrodynamic models: Coherent vibrations nucleonic fluids
(well described with liquid drop model(s))

Giant Resonances

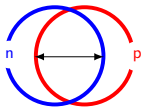
Hydrodynamic models

monopole
 $\Delta L = 0$
 $(\hat{\mathcal{O}} = Y_{0,0})$

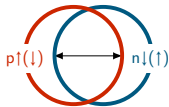
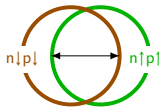


dipole
 $\Delta L = 1$
 $(\hat{\mathcal{O}} = Y_{1,\mu})$

ISGDR



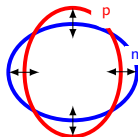
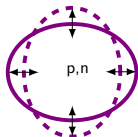
IVGDR



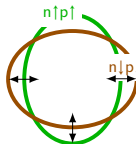
quadrupole
 $\Delta L = 2$
 $(\hat{\mathcal{O}} = Y_{2,\mu})$

ISGQR

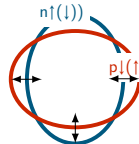
$\Delta T = 0$
 $\Delta S = 0$



$\Delta T = 1$
 $\Delta S = 0$
 $(\hat{\mathcal{O}} = \hat{t}_z)$



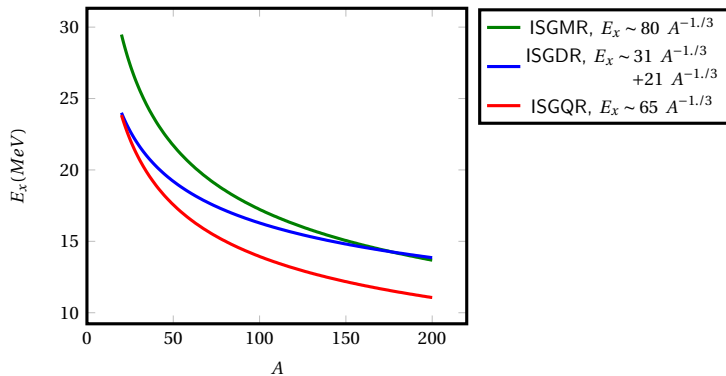
$\Delta T = 0$
 $\Delta S = 1$
 $(\hat{\mathcal{O}} = \hat{\sigma}_z)$



$\Delta T = 1$
 $\Delta S = 1$
 $(\hat{\mathcal{O}} = \hat{t}_z, \hat{\sigma}_z)$

Giant Resonances

Energies

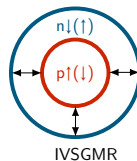
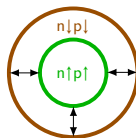
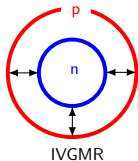
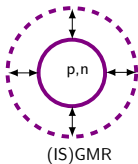


Harakeh et al. (2001)

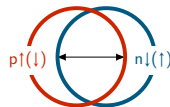
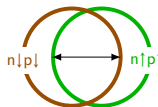
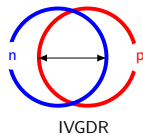
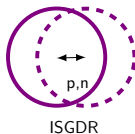
Giant Resonances

Hydrodynamic models

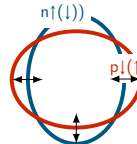
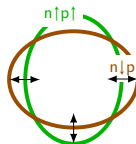
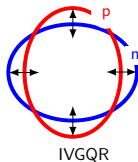
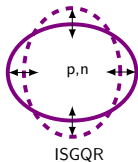
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 $\Delta L = 0$
 $(\hat{\mathcal{O}} = Y_{0,0})$



dipole
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quadrupole
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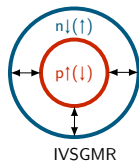
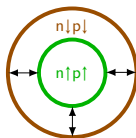
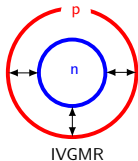
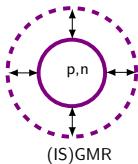
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Giant Resonances

Hydrodynamic models

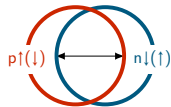
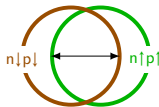
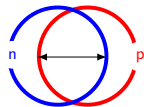
monopole
 $\Delta L = 0$
 $(\hat{\mathcal{G}} = Y_{0,0})$



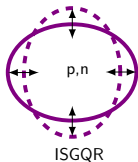
dipole
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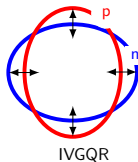
ISGDR



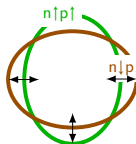
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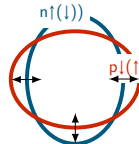
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Giant Resonances

Compression mode

Giant Resonances

Macroscopic/Hydrodynamic models: Coherent vibrations nucleonic fluids

Compression modes: ISGMR, ISGDR

Giant Monopole Resonance

IsoScalar GDR

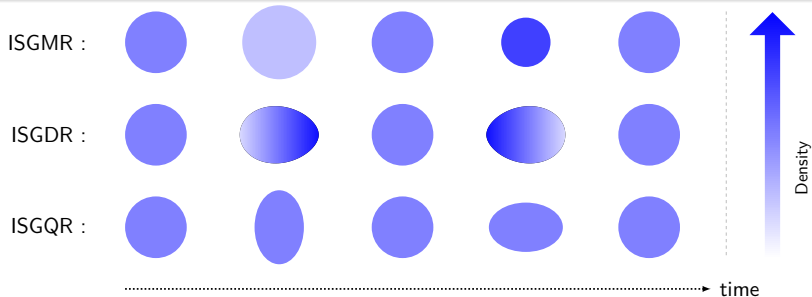
Giant Resonances

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Giant Resonances

Compression mode

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$$E_{\text{ISGMR}} = \hbar \sqrt{\frac{K_A}{m \langle r^2 \rangle}}$$

Stringari Phys. Lett. B 108 (1982)

Bagchi PhD (2015)

$$E_{\text{ISGDR}} = \hbar \sqrt{\frac{7}{3} \frac{K_A + \frac{27}{25} \epsilon_F}{m \langle r^2 \rangle}}$$

Nuclear incompressibility

$$K_A(\text{finite matter}) \neq K_\infty = \left[9\rho_0^2 \frac{\partial^2(\mathcal{E}/A)}{\partial \rho^2} \right]_{\rho_0} (\text{infinite matter})$$

ISGMR, ISGDR \Rightarrow Incompressibility + symmetry energy

Giant Resonances

Compression mode

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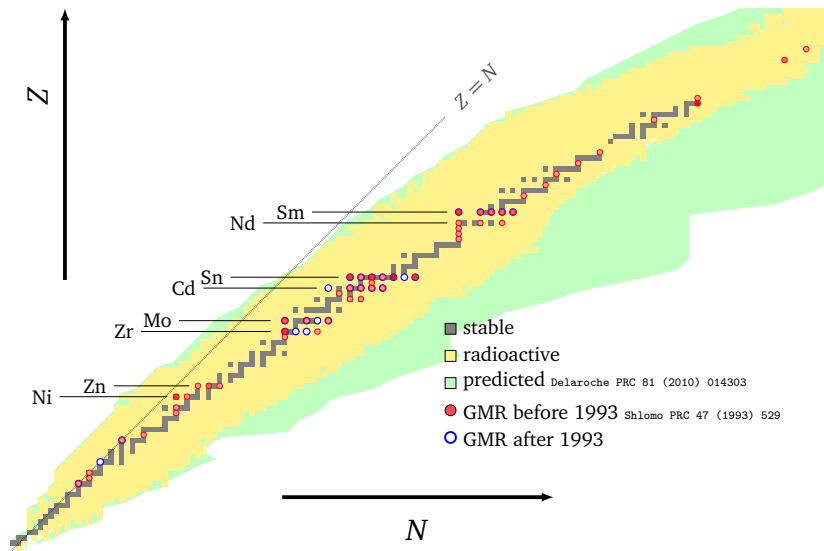
Nuclear incompressibility

$$K_A = K_\infty + K_S A^{-1/3} + K_T \underbrace{\left[\frac{N-Z}{A} \right]^2}_{\delta^2} + K_C \left[\frac{Z}{A^{1/3}} \right]^2 \quad (\text{scaling model})$$

ISGMR, ISGDR \Rightarrow Incompressibility + symmetry energy

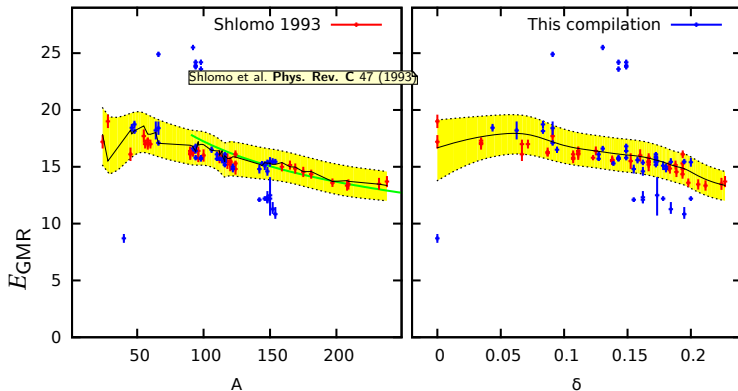
Giant Monopole Resonances

Status: stables nuclei only



Giant Monopole Resonances

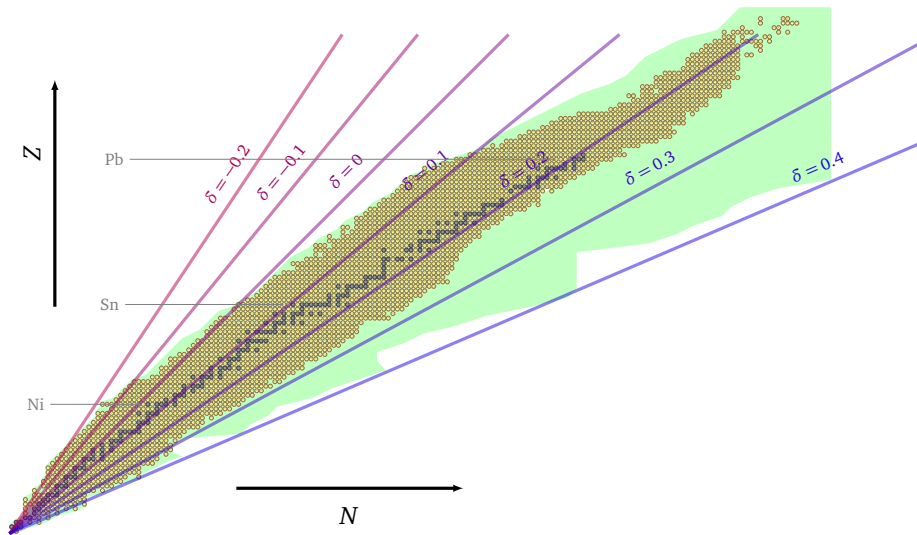
Incompressibility parameters



	This work			Shlomo
K_{∞}	220	240	260	244 ± 15
K_T	-28 ± 104	-222 ± 166	-416 ± 112	-142 ± 133

Giant Monopole Resonances

Towards exotic nuclei



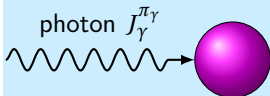
Experimental considerations

Properties of the probes

Selection rules

Historically GRs discovered by photon absorption (IVGDR)

Electromagnetic probes ?



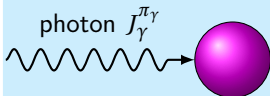
$$\Delta J = J_\gamma > 0 \text{ and}$$
$$\pi_\gamma = (-1)^{J_\gamma} \text{ (E) or } (-1)^{J_\gamma+1} \text{ (M)}$$

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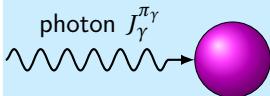
GMR: $\Delta L = \Delta S = 0 \Rightarrow \Delta J = 0$ EM \Leftrightarrow (virtual) photon

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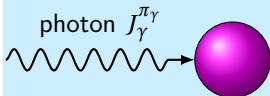
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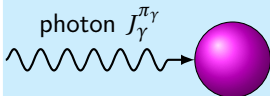
GDR: $\Delta L = 1$ & $\Delta S = 0 \Rightarrow \Delta J = 1$ & $\Delta\pi = (-1)^{\Delta L} = -1 \Rightarrow E1$

Properties of the probes

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SGDR: $\Delta L = 1, \Delta S = 1, \Delta\pi = -1 \Rightarrow \Delta J = 0, 1, 2 \Rightarrow E1$ or $M2$

Properties of the probes

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Inelastic scattering (α, α')

$$\left\{ \begin{array}{l} \vec{T}_i + \Delta \vec{T}_{GR} = \vec{T}_f \\ T_{z_i} + \Delta T_{z_{GR}} = T_{z_f} \end{array} \right. \quad \text{and} \quad \left\{ \begin{array}{l} \vec{T}_i + \vec{T}_\alpha = \vec{T}_f + \vec{T}_{\alpha'} \\ T_{z_i} + \cancel{T_{z_\alpha}}^0 = T_{z_f} + \cancel{T_{z_{\alpha'}}}^0 \end{array} \right.$$

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$$\Rightarrow \Delta\vec{T}_{GR} = \vec{T}_f - \vec{T}_i = \vec{T}_\alpha - \vec{T}_{\alpha'} \Rightarrow |T_\alpha - T_{\alpha'}| \leq \Delta T_{GR} \leq T_\alpha + T_{\alpha'} \Rightarrow 0 \leq \Delta T_{GR} \leq 2T_\alpha$$

Properties of the probes

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$\alpha(T=0)$: perfect probe for isoscalar modes (+ $S_{n,p,d,\dots} \gtrsim 20\text{MeV}$)

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Inelastic scattering

α ($T = 0$): perfect probe for isoscalar modes (+ $S_{n,p,d...} \gtrsim 20$ MeV)

d ($T = 0$): good probe for isoscalar modes also (but breaks! $S_n = S_p \approx 2$ MeV)

Properties of the probes

Selection rules

Historically GRs discovered by photon absorption (IVGDR)

Electromagnetic probes ?

GMR: ~~$\Delta L = \Delta S = 0 \Rightarrow \Delta J = 0$~~ EM \leftrightarrow (virtual) photon

GDR: $\Delta L = 1$ & $\Delta S = 0 \Rightarrow \Delta J = 1$ & $\Delta\pi = (-1)^{\Delta L} = -1 \Rightarrow E1$

SGDR: $\Delta L = 1, \Delta S = 1, \Delta\pi = -1 \Rightarrow \Delta J = 0, 1, 2 \Rightarrow E1$ or $M2$

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More general probes

	$\Delta S = 0$ $\Delta T = 0$ $\Delta A = 0$	$\Delta S = 1$ $\Delta T = 0$ $\Delta A = 0$	$\Delta S = 0$ $\Delta T = 1$ $\Delta A = 0$	$\Delta S = 1$ $\Delta T = 1$ $\Delta A = 0$	ΔS ΔT $\Delta A = 2$
Variable	Number density	Spin density	Isvector density	Isvector spin density	Pair density
Property	Incompressibility	Magnetism	Symmetry energy		Pair condensation
Probe	(α, α) , (d, d)	(p, p) , $({}^6\text{Li}, {}^6\text{Li}^*)$	$({}^7\text{Li}, {}^7\text{Be}^*)$, $({}^6\text{He}, {}^6\text{Li}^*)$	(p, n) , (n, p) , $(d, 2p)$...	$(\alpha, {}^6\text{He})$, $(\alpha, {}^6\text{Li})$, (α, d) , (α, pn) , (d, α) , $(n, {}^3\text{He})$, $({}^3\text{He}, n)$

Properties of the probes

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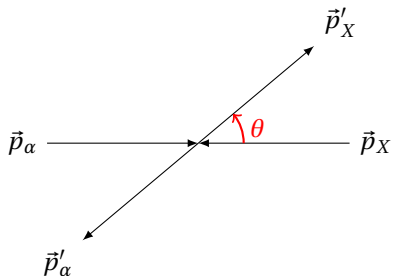
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Angular distribution : General considerations

Classical approach

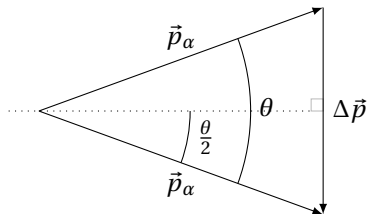
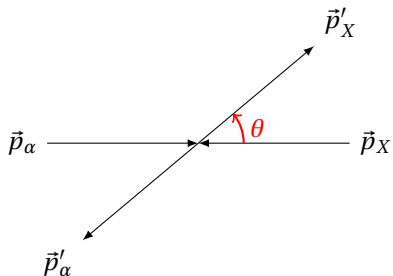
$${}^A X(\alpha, \alpha') {}^A X^*$$



Angular distribution : General considerations

Classical approach

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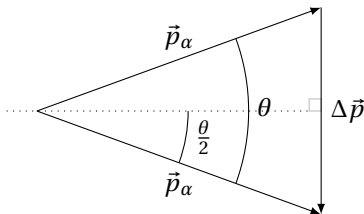
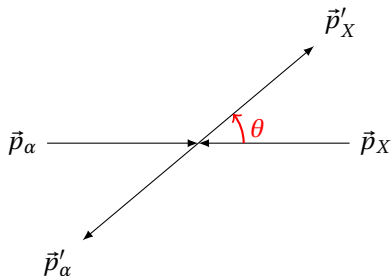


The momentum transferred by α is : $\Delta \vec{p} = \vec{p}'_{\alpha} - \vec{p}_{\alpha}$.

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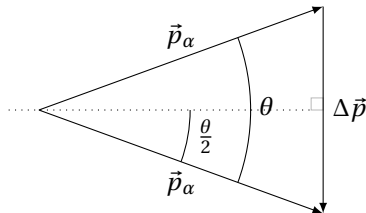
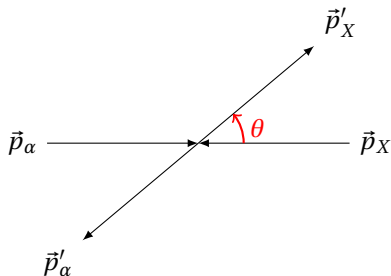
If we assume $p_{\alpha} \sim p'_{\alpha}$: $\sin \frac{\theta}{2} \sim \frac{|\Delta \vec{p}|/2}{|\vec{p}_{\alpha}|}$

with $\vec{l}_{\alpha} = \vec{r} \wedge \Delta \vec{p} \Rightarrow l_{\alpha} = r \Delta p \Rightarrow \Delta p = l_{\alpha} / r \Rightarrow \sin(\theta/2) = \frac{l_{\alpha}}{2p_{\alpha}r} \Rightarrow \theta = 2 \arcsin \left(\frac{|\vec{l}_{\alpha}|}{2|\vec{p}_{\alpha}||\vec{r}|} \right)$

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$$\theta \propto l$$

Angular distribution shape

With interferences

Sharp edge Fraunhofer (approximation w/o Coulomb)

Garg et al. *Progress in Particle and Nuclear Physics* 101 (2018) Bernstein (1969)

(Analogy w/ optics)

$$\begin{aligned}\left(\frac{d\sigma}{d\Omega}\right)_{0^+ \rightarrow 0^+} &\propto |J_0(qR_D)|^2, \\ \left(\frac{d\sigma}{d\Omega}\right)_{0^+ \rightarrow 1^-} &\propto |J_1(qR_D)|^2, \\ \left(\frac{d\sigma}{d\Omega}\right)_{0^+ \rightarrow 2^+} &\propto \left[\frac{1}{4}J_0(qR_D)^2 + \frac{3}{4}J_2(qR_D)^2\right] \propto J_2^2(qR_D)\end{aligned}$$

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Angular distribution shape

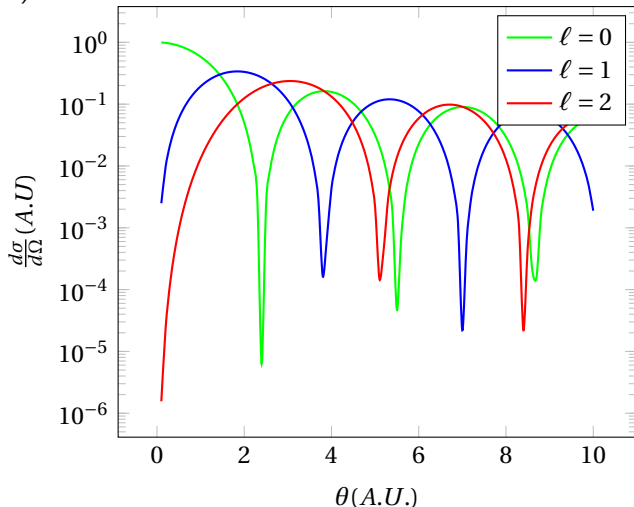
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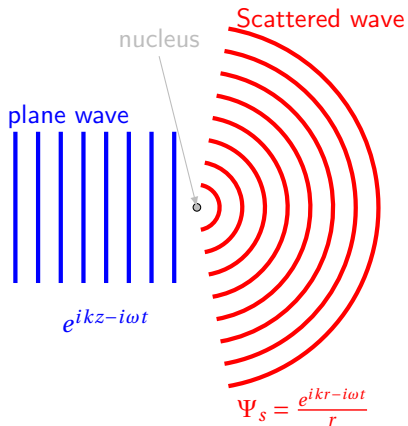
(Analogy w/ optics)

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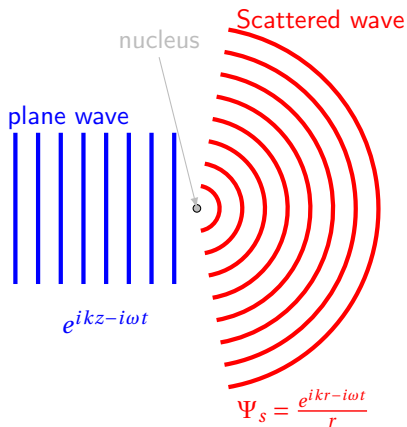
Angular distribution calculation

In practice: optical potential, DW ...



Angular distribution calculation

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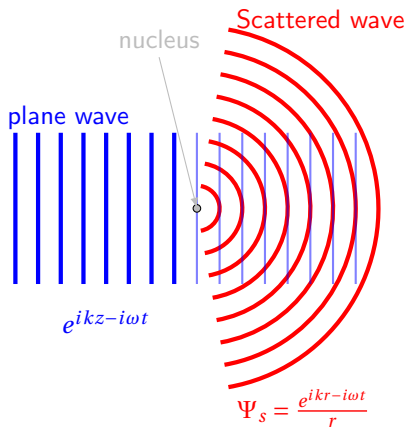


Total wave function:

$$\psi_{\text{tot}}(\theta) \xrightarrow{r \rightarrow \infty} \left[e^{ikz} + \underbrace{f(\theta)}_{\psi_s} \frac{e^{ikr}}{r} \right] e^{-i\omega t}$$

Angular distribution calculation

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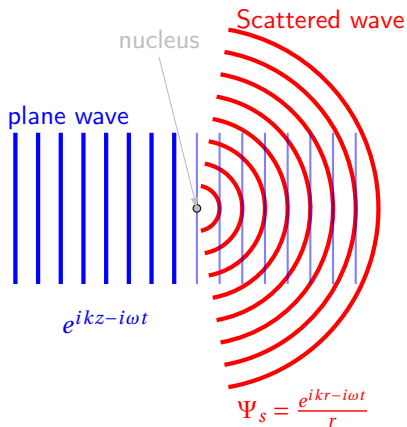


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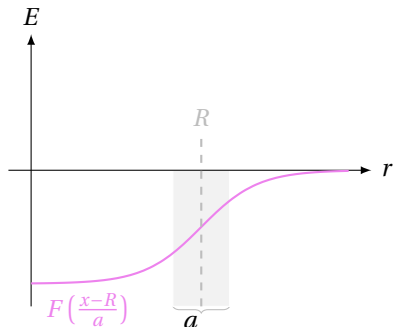
Angular distribution

$$\frac{d\sigma}{d\Omega} = |f(\theta)|^2$$

Angular distribution calculation

In practice: optical potential, DW ...

Nucleus representation:

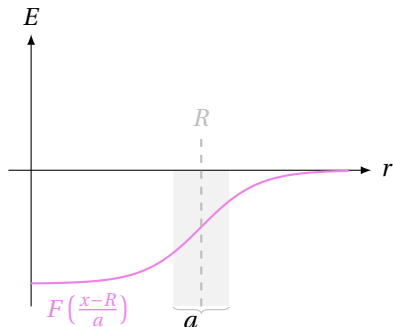


Wood saxon potential
(w/o Coulomb)

Angular distribution calculation

In practice: optical potential, DW ...

Nucleus representation:



Wood saxon potential
(w/o Coulomb)

To calculate ψ :

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + \mathbf{V}(r) \right] \psi = E\psi$$

with *ad minima*

$$\mathbf{V}(r) = -(\mathcal{V} + i\mathcal{W}) F\left(\frac{r-R}{a}\right)$$

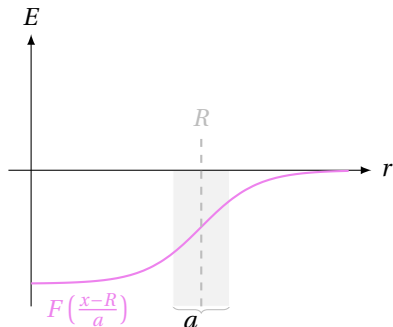
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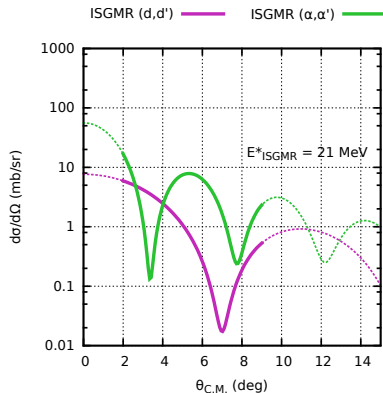
$$\psi = \sum_l \sum_{m=-l}^l \alpha_{l,m}(k) \frac{u_l(k,r)}{r} Y_l^m(\theta, \psi)$$

+ deformation/rotation + form factor (Harakeh et al. *Physical Review C* 23 (1981)) ...

Giant Resonances

Experimental techniques?

$^{68}\text{Ni}(\alpha, \alpha')^{68}\text{Ni}^*$

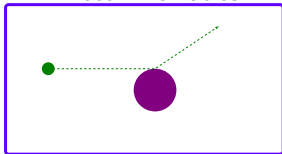


Physics of interest $\sim 0^\circ$

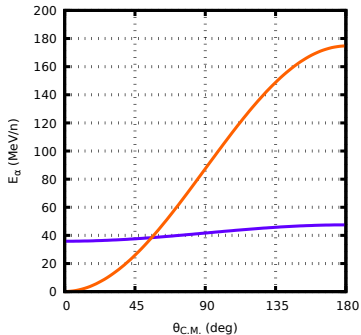
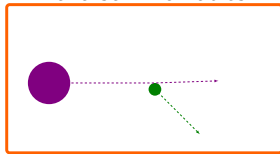
Kinematics

Example : α on $A \sim 60$ with $T \sim 50 \text{ MeV/n}$

Direct kinematics



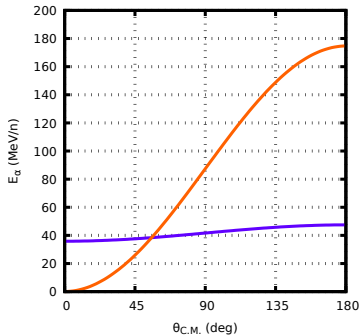
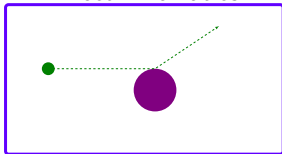
Reverse kinematics



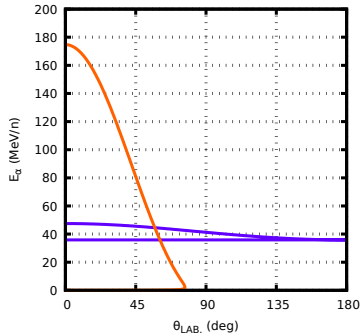
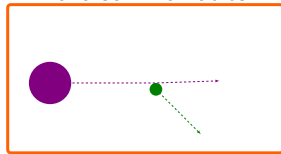
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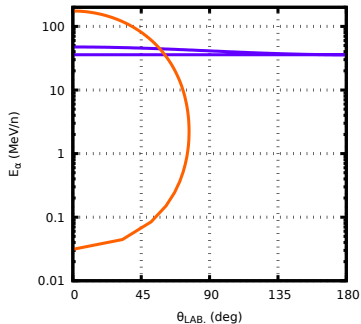
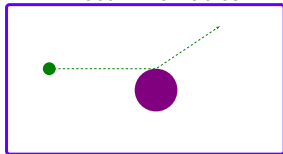
Reverse kinematics



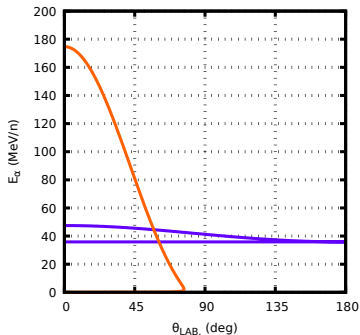
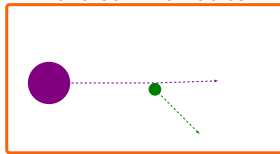
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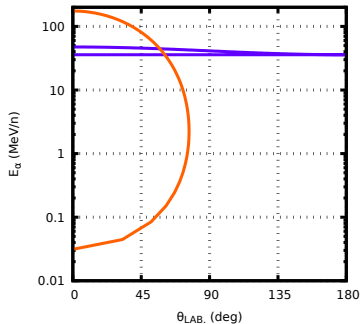
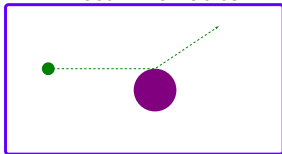
Reverse kinematics



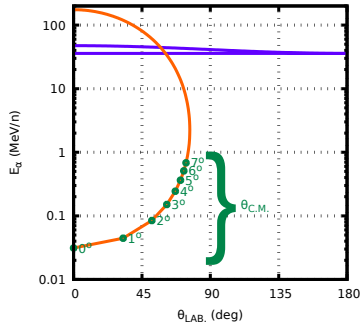
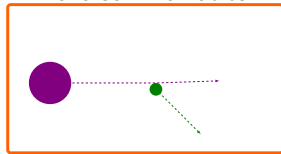
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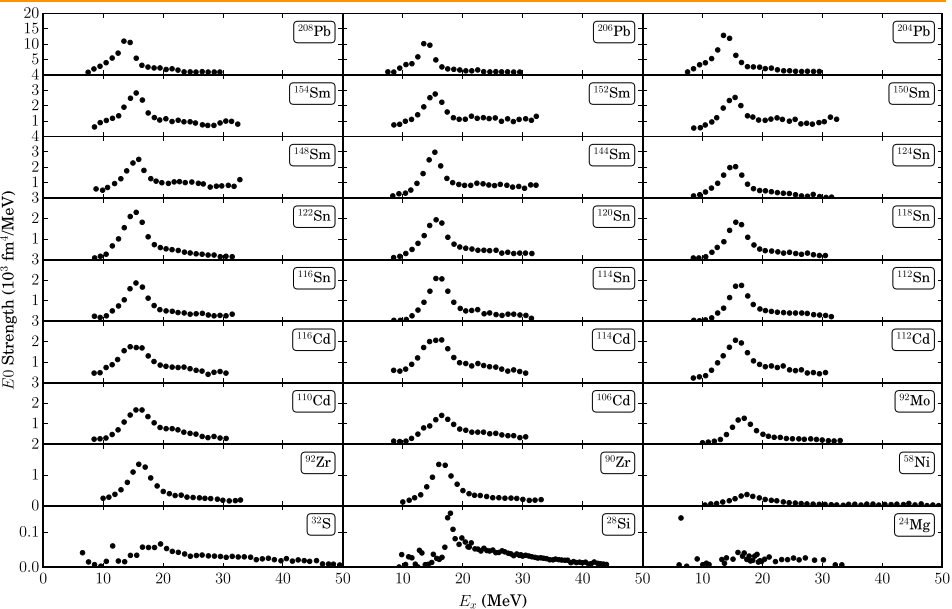
Direct kinematics

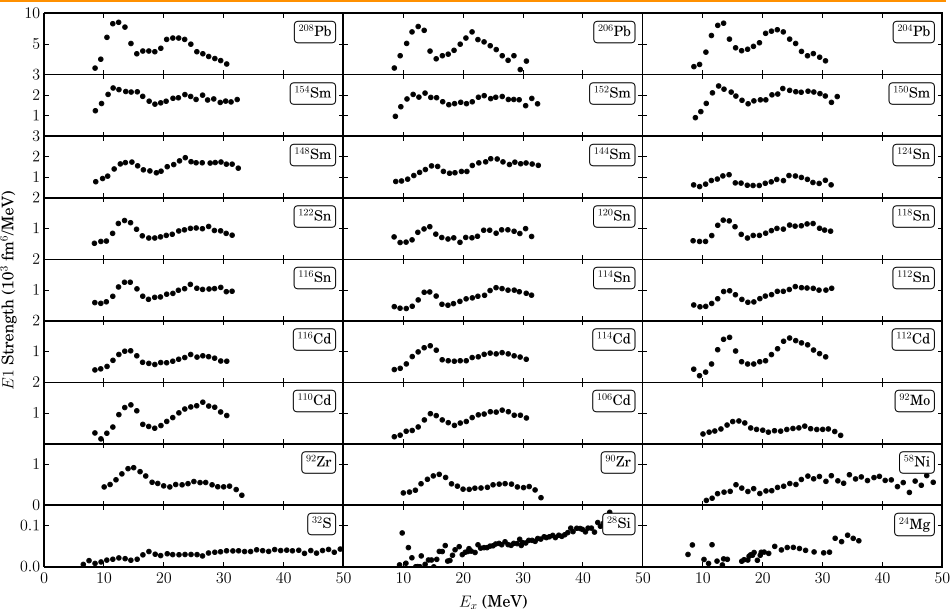


Reverse kinematics



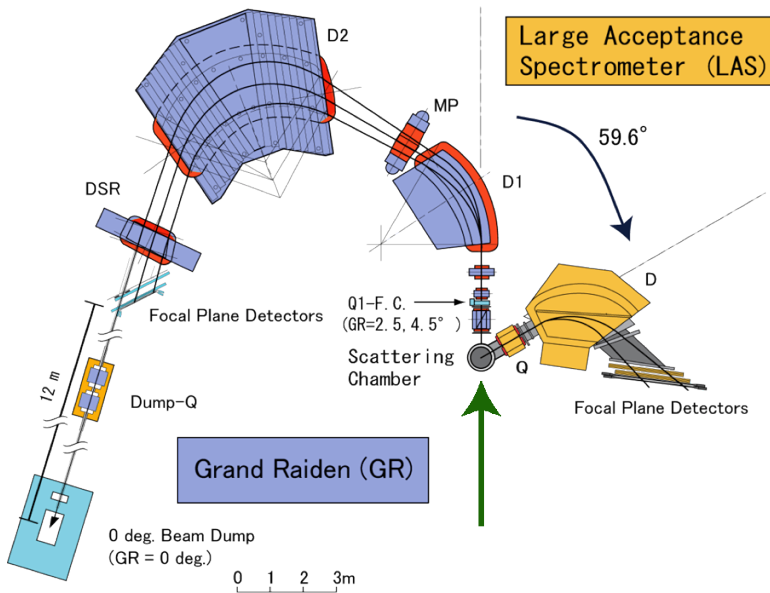
Example of experiments with stable nuclei





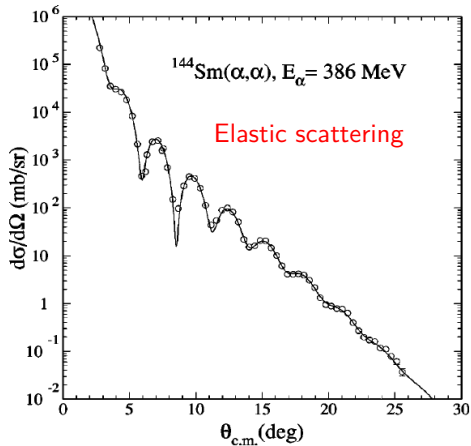
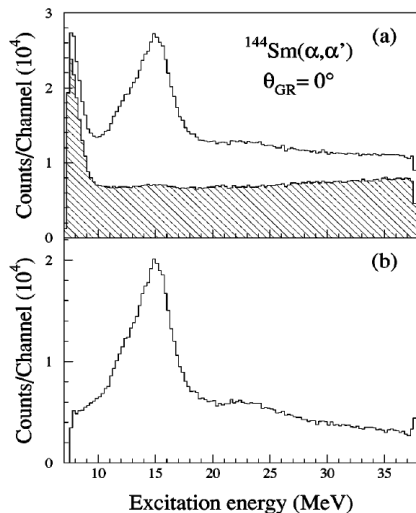
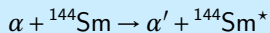
Systematics on stable nuclei

Experimental details [Itoh et al. Phys. Rev. C 68 \(2003\)](#)



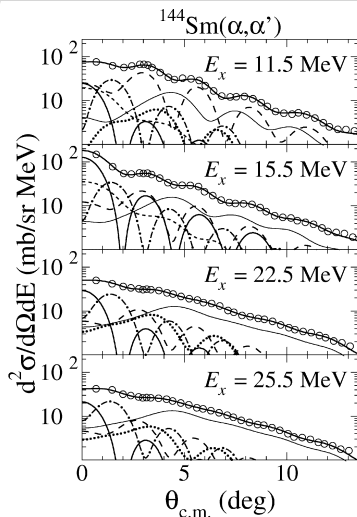
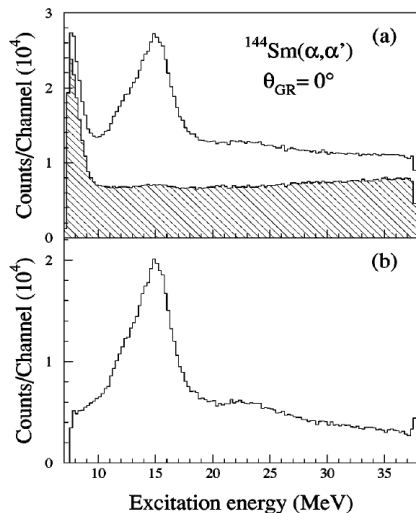
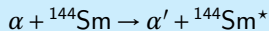
Systematics on stable nuclei

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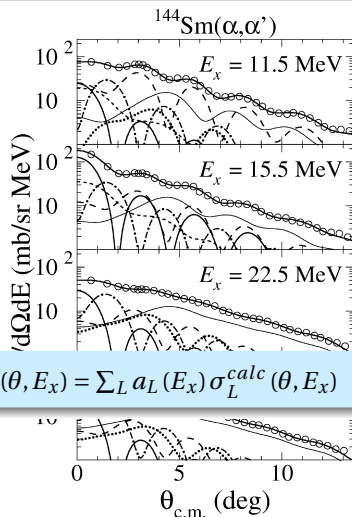
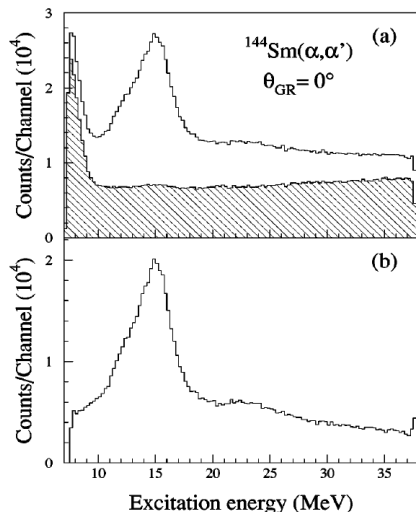
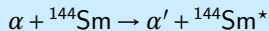
Systematics on stable nuclei

Experimental details Itoh et al. Phys. Rev. C 68 (2003)



Systematics on stable nuclei

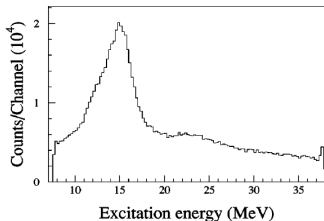
Experimental details Itoh et al. Phys. Rev. C 68 (2003)



$$\sigma^{exp}(\theta, E_x) = \sum_L a_L(E_x) \sigma_L^{calc}(\theta, E_x)$$

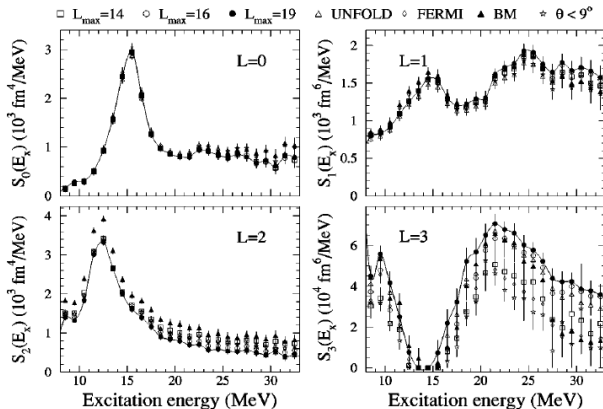
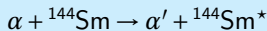
Systematics on stable nuclei

Analysis details Itoh et al. Phys. Rev. C 68 (2003)



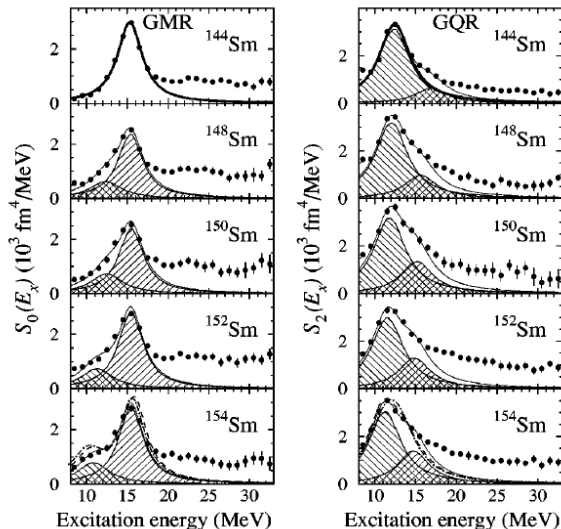
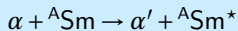
Multiple
Decomposition
Analysis
(MDA)

A red arrow points from the experimental spectrum in the top-left figure to the decomposed components in the bottom-right figure, indicating the application of Multiple Decomposition Analysis (MDA).



Systematics on stable nuclei

Analysis details Itoh et al. Phys. Rev. C 68 (2003)



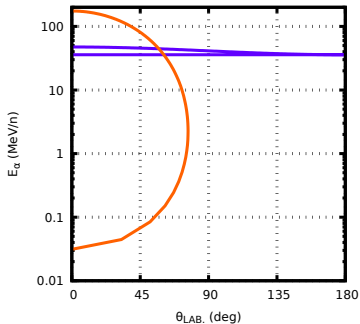
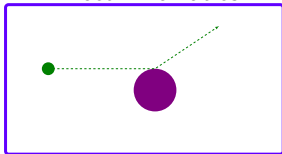
Finally fit
with Lorentzian
↓
position
width
“amplitude”
(X-section)

Example of experiments with unstable nuclei

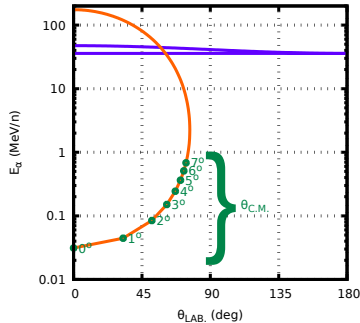
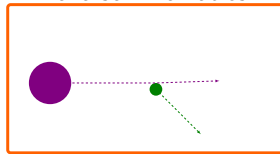
Kinematics

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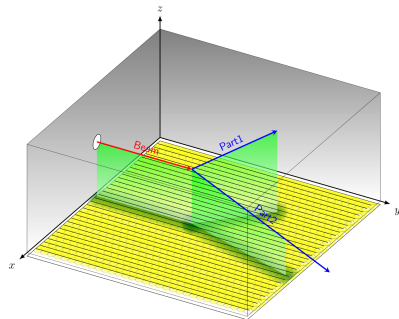
Direct kinematics



Reverse kinematics



Reverse kinematics \Rightarrow very low energy recoil \Rightarrow gas detector/target!



Gas detector w/ gas = target

Pros

High efficiency (3D)

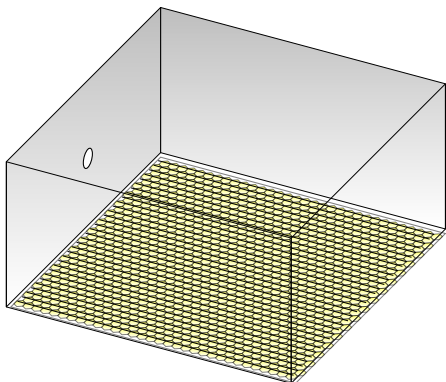
Low detection threshold

Thick target ($\approx 10 \text{ mg/cm}^2$)

Cons (as for MAYA)

Beam rate limited (10^5 pps)

Capricious (!)



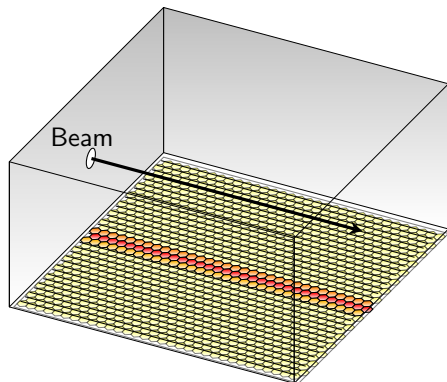
MAYA

$283 \times 258 \times 200 \text{ mm}^3 \equiv 30 \text{ L of gas}$

⇒ thick target

⇒ gas detector

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)



MAYA

$283 \times 258 \times 200 \text{ mm}^3 \equiv 30 \text{ L of gas}$

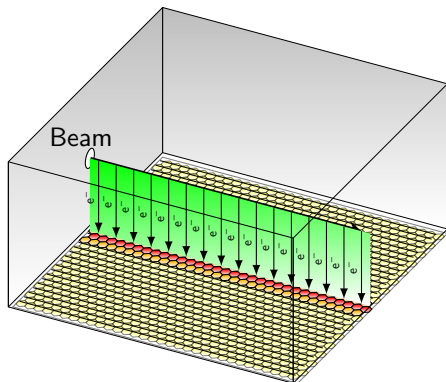
⇒ thick target

⇒ gas detector

Inside :

2D proj. (honeycomb struct.)

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)



MAYA

$283 \times 258 \times 200 \text{ mm}^3 \equiv 30 \text{ L of gas}$

⇒ thick target

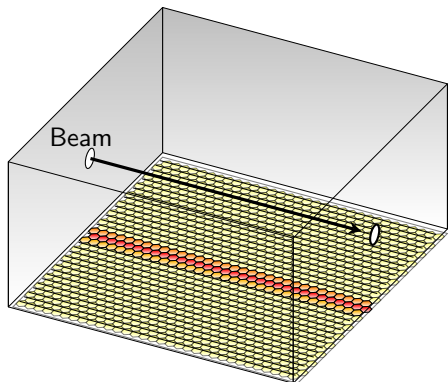
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Inside :

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3D using drift time

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)



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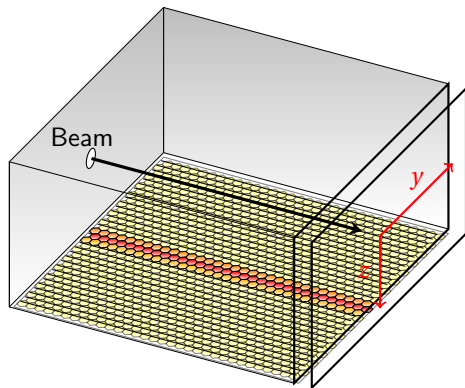
2D proj. (honeycomb struct.)

3D using drift time

Ancillary detectors :

Diamond

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)



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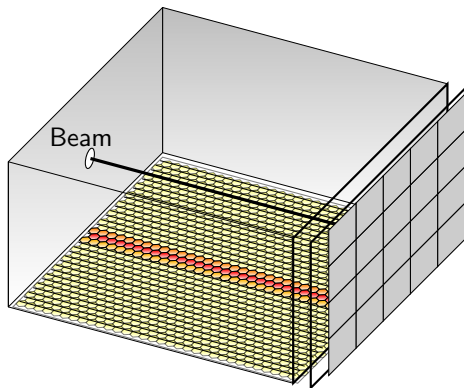
3D using drift time

Ancillary detectors :

Diamond

Drift chamber

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)



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3D using drift time

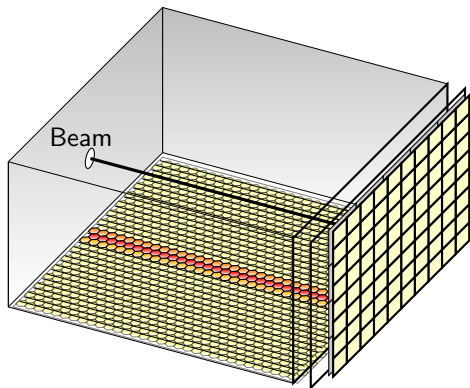
Ancillary detectors :

Diamond

Drift chamber

4×5 Si ($5 \times 5 \text{ cm}^2$, $700 \mu\text{m}$)

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)



MAYA

$283 \times 258 \times 200 \text{ mm}^3 \equiv 30 \text{ L of gas}$

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2D proj. (honeycomb struct.)

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Diamond

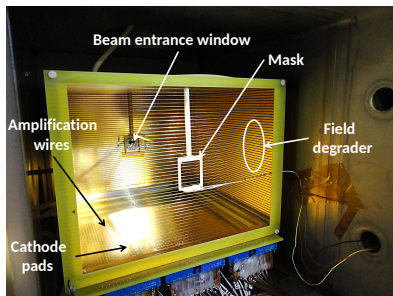
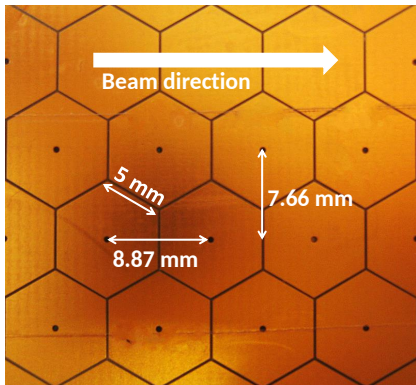
Drift chamber

$4 \times 5 \text{ Si } (5 \times 5 \text{ cm}^2, 700 \text{ } \mu\text{m})$

$8 \times 10 \text{ CsI } (2.5 \times 2.5 \text{ cm}^2)$

Demonchy et al. Nucl. Instrum. Methods Phys. Res., Sect. A 583 (2007)

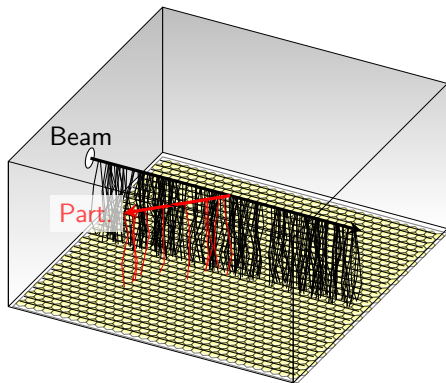
Setup: pictures



Setup: Improvement

Mask

Beam energy loss \gg scattered particles \Rightarrow dynamics issues.

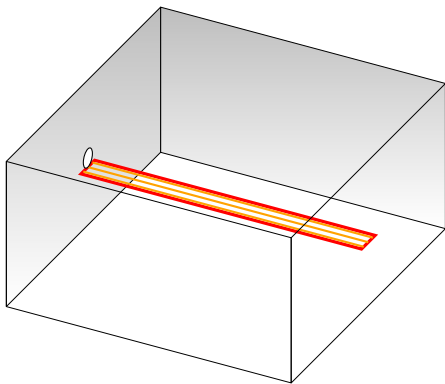


Pancin et al. *J. Instrum.* 7 (2012)

Setup: Improvement

Mask

Beam energy loss \gg scattered particles \Rightarrow dynamics issues.

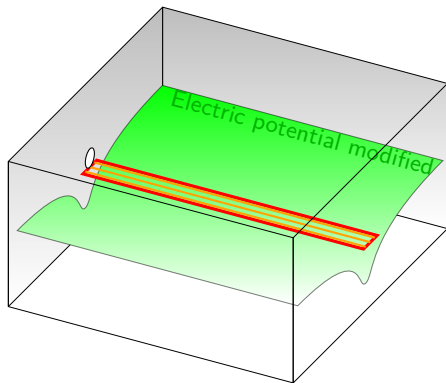


Pancin et al. *J. Instrum.* 7 (2012)

Setup: Improvement

Mask

Beam energy loss \gg scattered particles \Rightarrow dynamics issues.

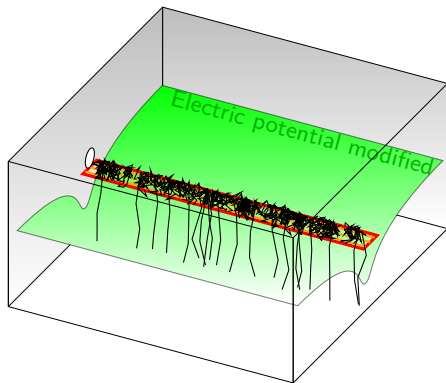


Pancin et al. *J. Instrum.* 7 (2012)

Setup: Improvement

Mask

Beam energy loss \gg scattered particles \Rightarrow dynamics issues.

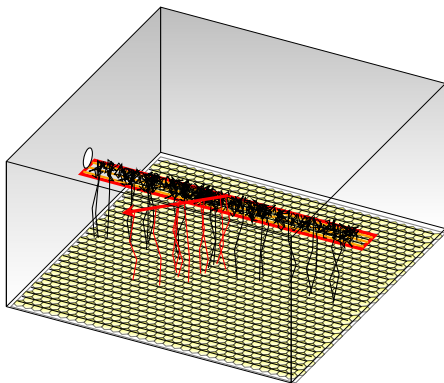


Pancin et al. *J. Instrum.* 7 (2012)

Setup: Improvement

Mask

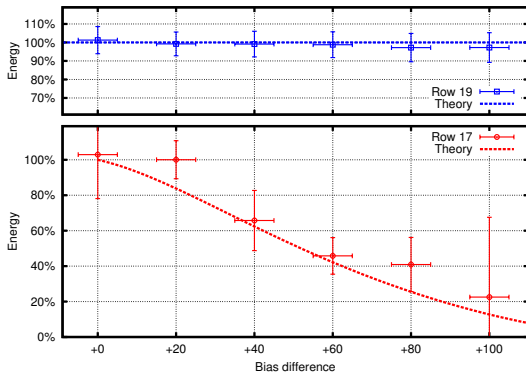
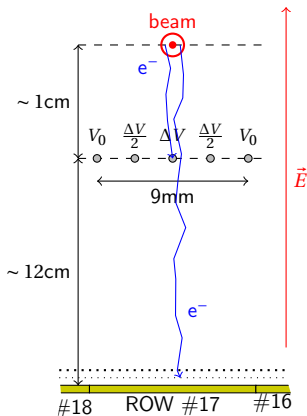
Beam energy loss \gg scattered particles \Rightarrow dynamics issues.



Pancin et al. *J. Instrum.* 7 (2012)

Setup: Improvement

Mask



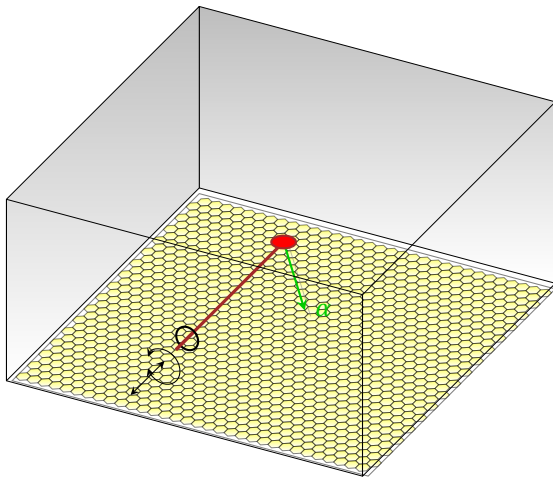
Bias difference (ΔV) added to **mask** central wire ; Energy $\propto N_{e^-}$

Pancin et al. *J. Instrum.* 7 (2012)

Setup: Improvement

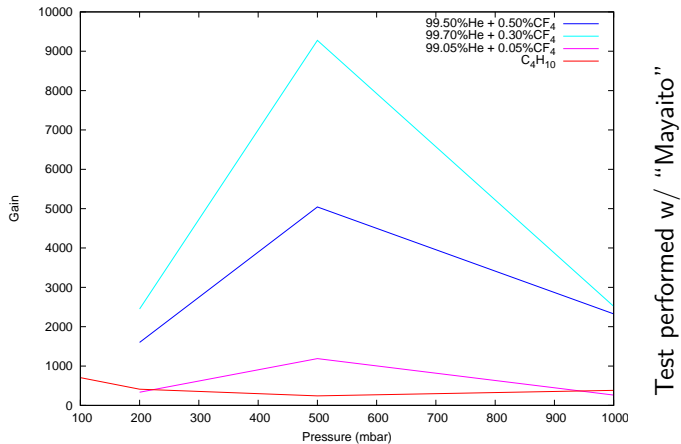
Source

Retractable source, allows regular monitoring.



Setup: Improvement

Gas: Helium feasible with CF_4 as quencher

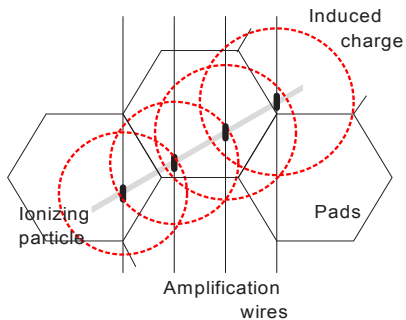
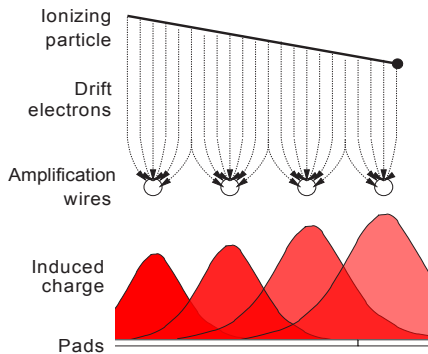


With MAYA and during experiment ~500mbar, 5% CF_4 .

Some MAYA analysis techniques:

Roger et al. Nucl. Instrum. Methods Phys. Res., Sect. A 638 (2011)

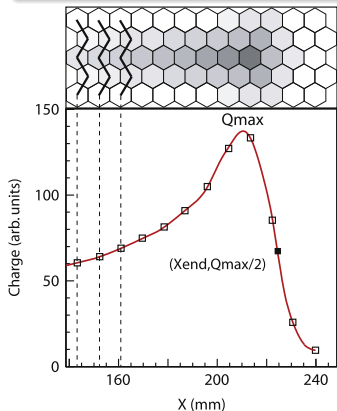
Trajectory reconstruction



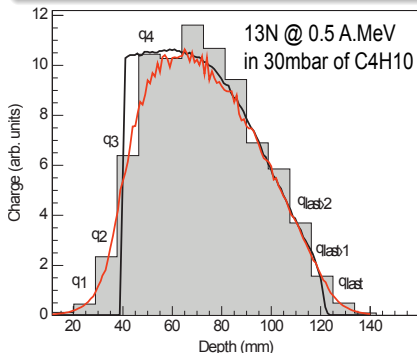
Some MAYA analysis techniques:

Range measurement

In case of a “clear” Bragg pic:
smoothing w/ spline \Rightarrow error ≤ 1 mm



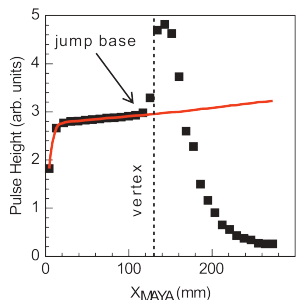
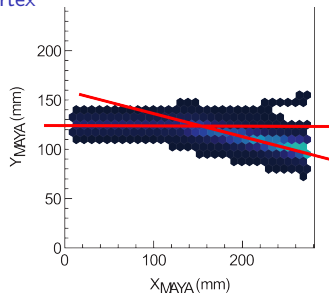
Otherwise: depends on the
experiment/simulation (typically last
charge/position combination...)



Roger et al. Nucl. Instrum. Methods Phys. Res., Sect. A 638 (2011)

Some MAYA analysis techniques:

Vertex



First order

intersection point of tracks

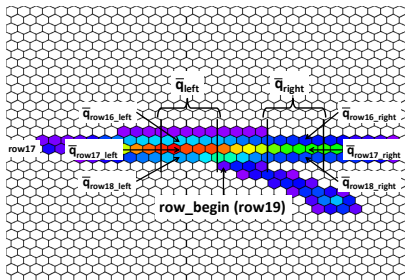
Refinement

enhancement of the deposited charge around the vertex. (simulation)

Some MAYA analysis techniques:

Vertex

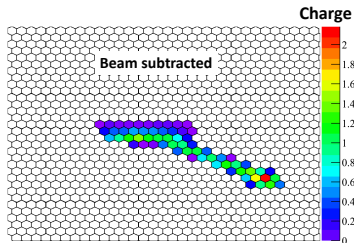
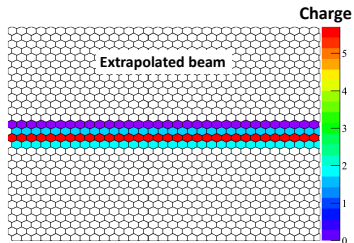
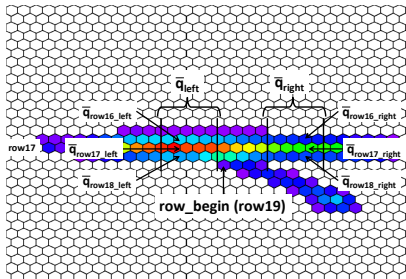
Thanks to the mask, beam can be seen and subtracted, to allow fitting of small tracks (few pads)



Some MAYA analysis techniques:

Vertex

Thanks to the mask, beam can be and subtracted, to allow fitting of small tracks (few pads)



Mounting MAYA !



<https://www.youtube.com/watch?v=aUxC68zb6sI>

^{56}Ni : ISGMR & ISGDR



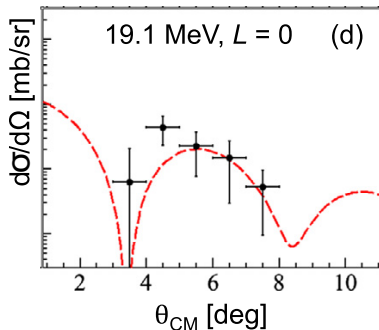
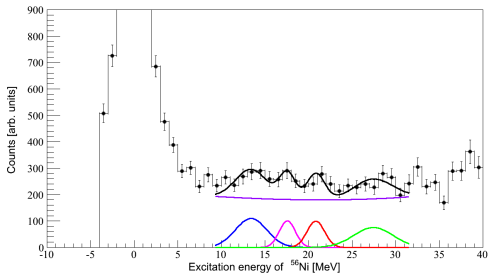
ISGMR : two experiments

$^{56}\text{Ni}(d, d')^{56}\text{Ni}^*$: [Monrozeau et al. Phys. Rev. Lett. 100 \(2008\)](#) [Monrozeau PhD \(2007\)](#)

$$E_{\text{ISGMR}}^* \approx 19.5 \text{ MeV}$$

$^{56}\text{Ni}(\alpha, \alpha')^{56}\text{Ni}^*$: [Bagchi PhD \(2015\)](#) [Bagchi et al. Phys. Lett. B 751 \(2015\)](#)

$$E_{\text{ISGDR}}^* = 19.1 \pm 0.5 \text{ MeV}$$



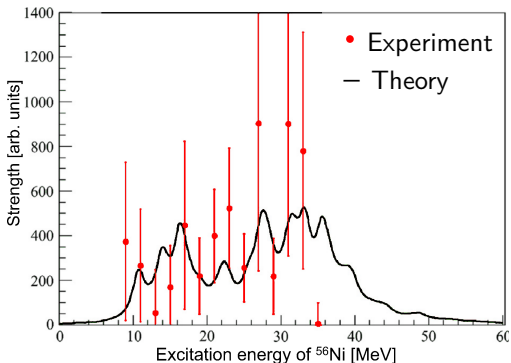
^{56}Ni : ISGMR & ISGDR



ISGDR

$^{56}\text{Ni}(\alpha, \alpha')^{56}\text{Ni}^*$: [Bagchi et al. Phys. Lett. B 751 \(2015\)](#)

$$E_{\text{ISGDR}}^* = 17.4 \pm 0.7 \text{ MeV}$$



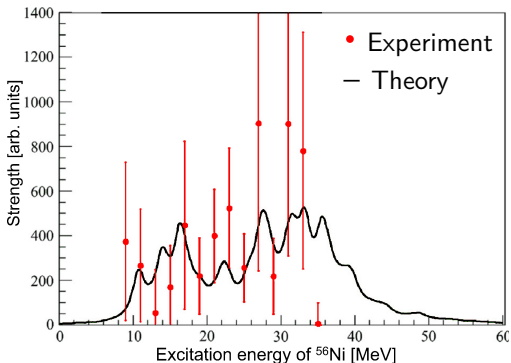
^{56}Ni : ISGMR & ISGDR



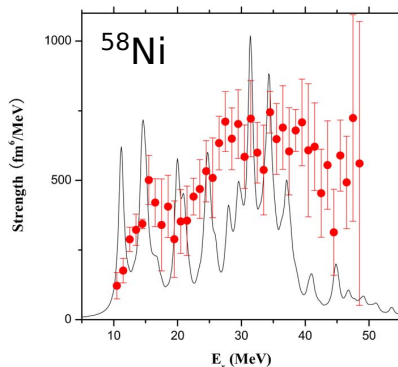
ISGDR

$^{56}\text{Ni}(\alpha, \alpha')^{56}\text{Ni}^*$: Bagchi et al. Phys. Lett. B 751 (2015)

$$E_{\text{ISGDR}}^* = 17.4 \pm 0.7 \text{ MeV}$$



This work



Nayak et al. Phys. Lett. B 637 (2006)

^{68}Ni : ISGMR



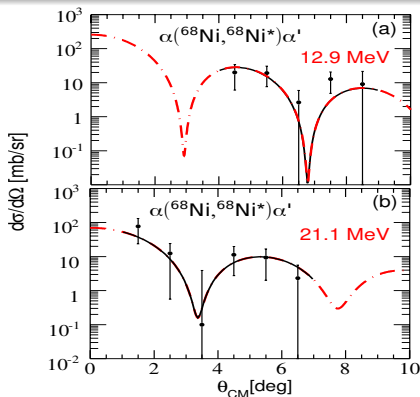
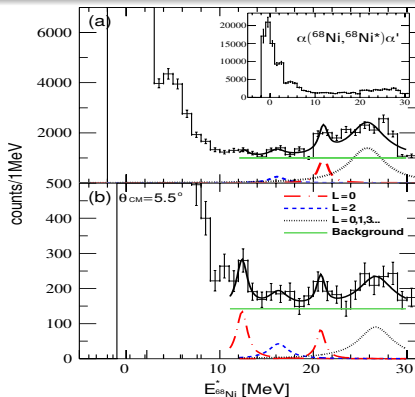
One proposal: two experiments (M. Vandebrouck's PhD)

$^{68}\text{Ni}(\alpha, \alpha')^{68}\text{Ni}^*$:

$E^* = 12.9 \pm 1.0$ & 21.1 ± 1.9 MeV

$^{68}\text{Ni}(d, d')^{68}\text{Ni}^*$:

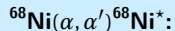
$E^* = 12.6 \pm 0.3$ & 20.8 ± 0.6 MeV



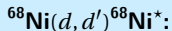
^{68}Ni : ISGMR



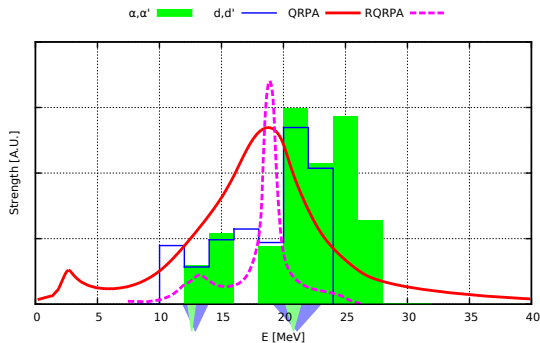
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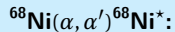
QRPA: [Terasaki et al. Phys. Rev. C 74 \(2006\)](#)

RQRPA: [Khan et al. Phys. Rev. C 84 \(2011\)](#)

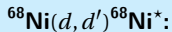
^{68}Ni : ISGMR



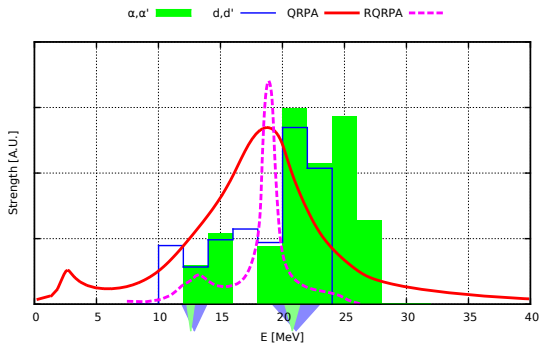
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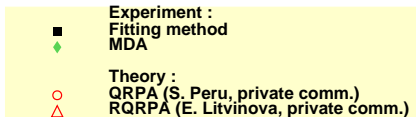
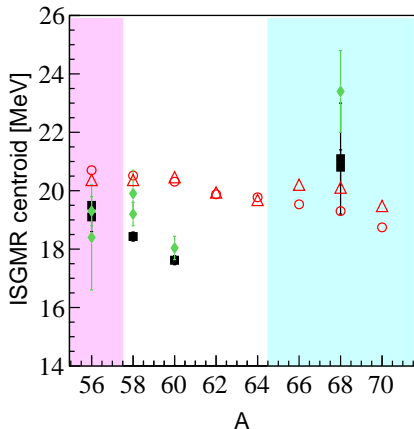


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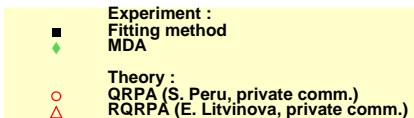
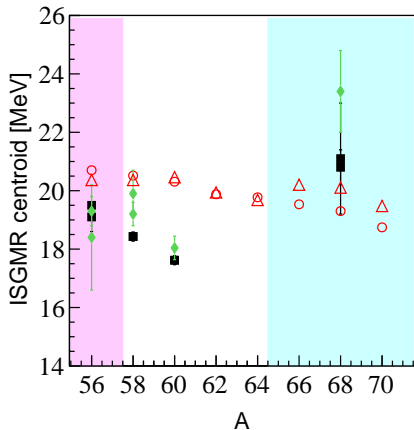


QRPA: [Terasaki et al. Phys. Rev. C 74 \(2006\)](#) QRPDA: [Khan et al. Phys. Rev. C 84 \(2011\)](#)

ISGMR : summary in Nickel's chain



ISGMR : summary in Nickel's chain



More points needed?

Improve measurement on $^{58,68,(70)}\text{Ni}$

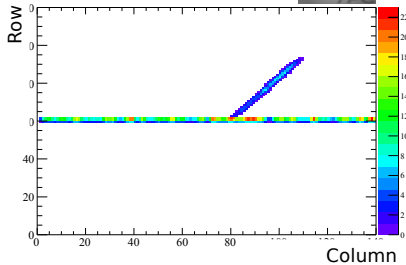
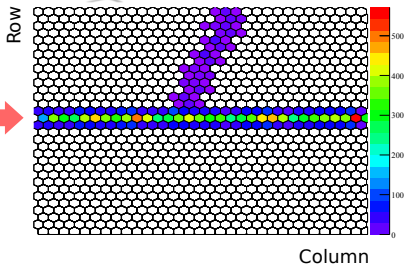


MAYA

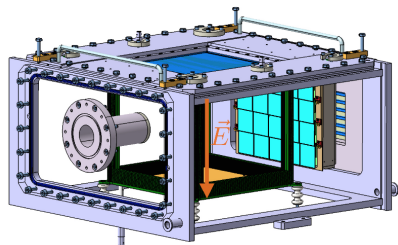
$^{68}\text{Ni}(\alpha, \alpha')^{68}\text{Ni}^*$ ($E^*=20\text{MeV}$)

actar
TPC

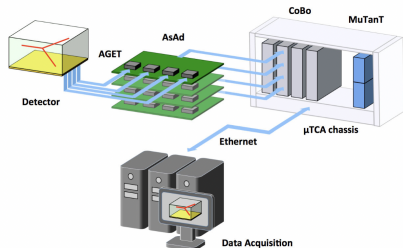
Beam
→



ACTAR TPC : ACTIVE TARget

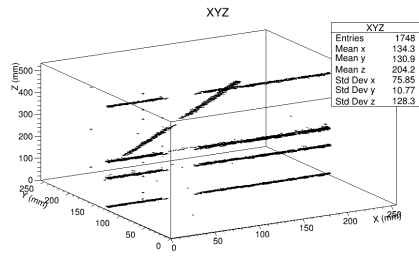
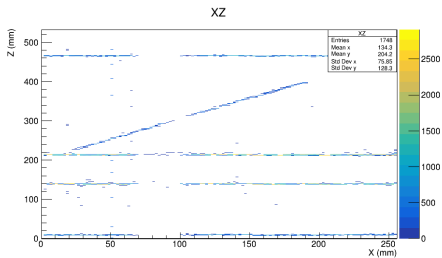
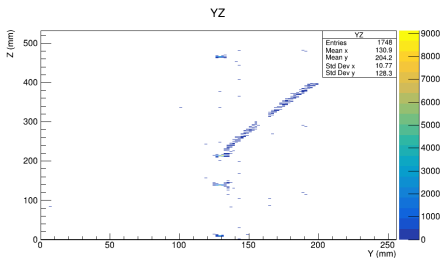
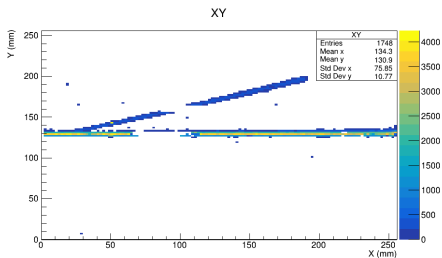


128×128 pixels of 2 mm^2
16384 channels in a square of
 $25.6 \times 25.6 \text{ cm}^2$ (~ 10 more
than MAYA)
Dedicated (compact) electronics



Roger et al. Nucl. Instrum. Methods Phys. Res., Sect. A 895 (2018)

ACTAR TPC : ACTIVE TARget



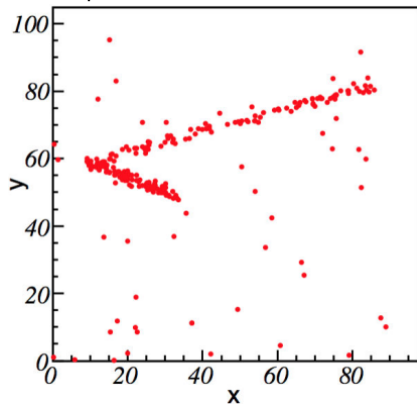
How to identify each individual track ?

ACTAR TPC : ACTive TARget

Identification and reconstruction of tracks

RANSAC method: iterative method to find the tracks

2D example

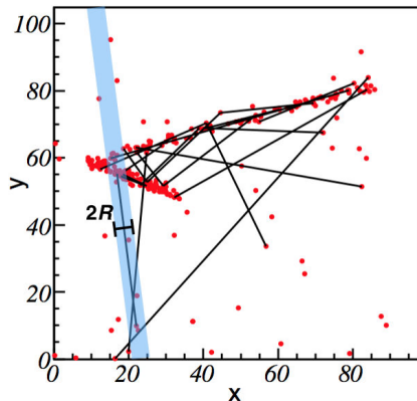


ACTAR TPC : ACTIVE TARget

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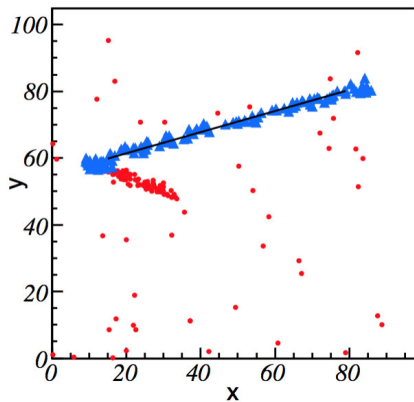
Mauss PhD (2018)

ACTAR TPC : ACTIVE TARget

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2D example



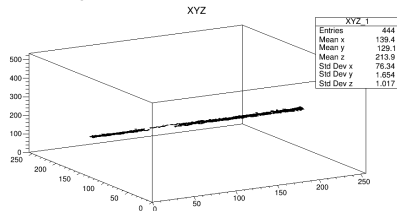
Mauss PhD (2018)

ACTAR TPC : ACTive TARget

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3D example iterative results

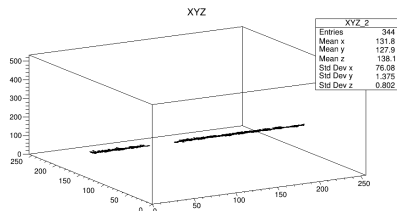
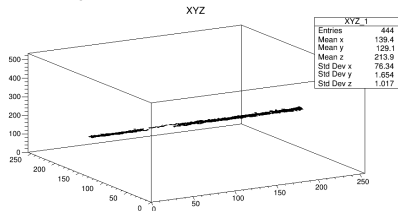


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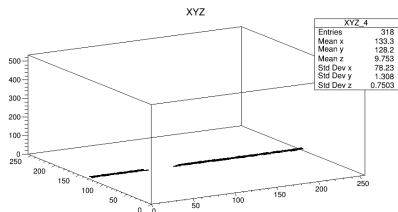
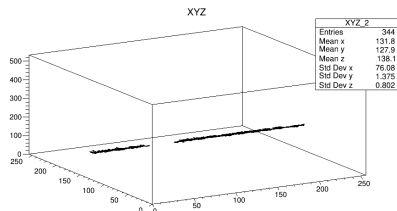
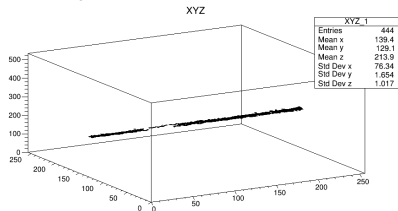


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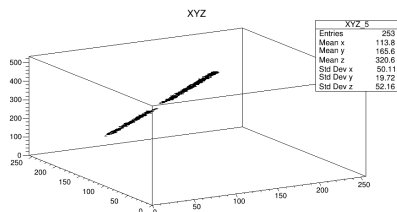
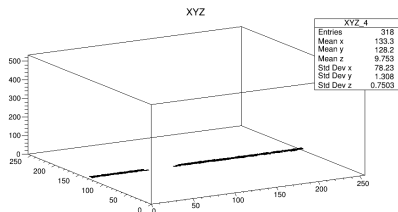
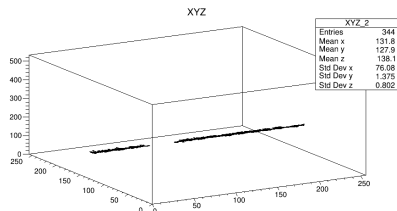
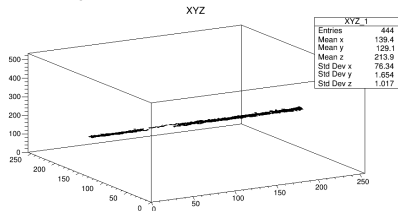
Mauss PhD (2018)

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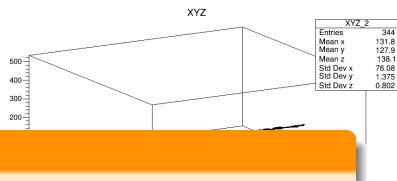
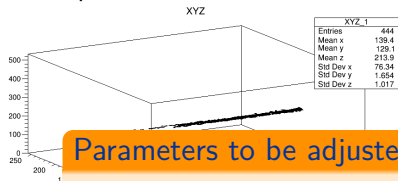
Mauss PhD (2018)

ACTAR TPC : ACTIVE TARget

Identification and reconstruction of tracks

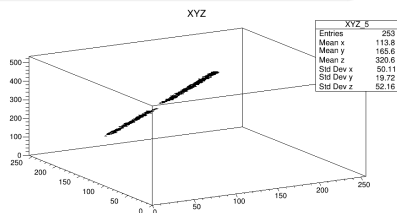
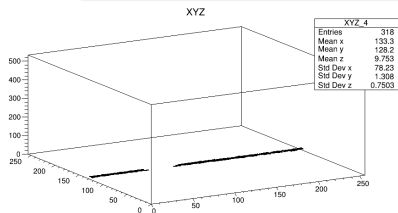
RANSAC method: iterative method to find the tracks

3D example iterative results



Parameters to be adjusted

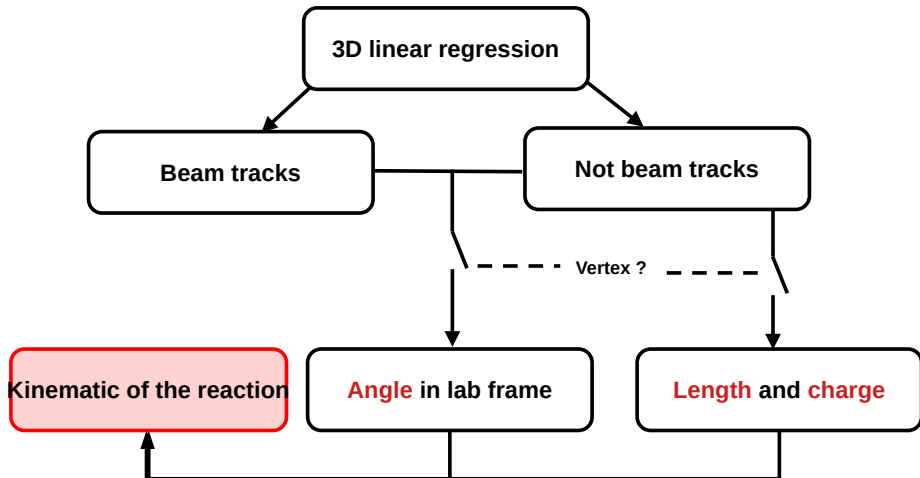
N and R ; Selection and stopping conditions



Mauss PhD (2018)

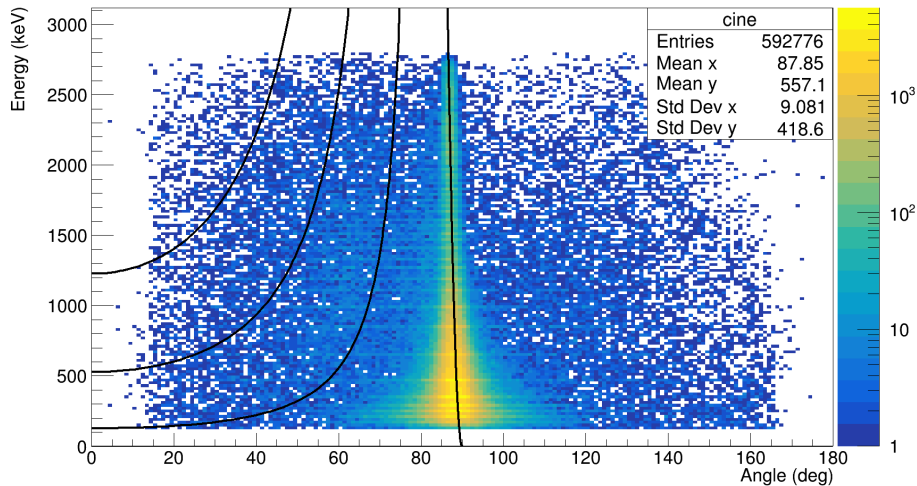
ACTAR TPC : ACTive TARget

Measurement of the kinematics parameters



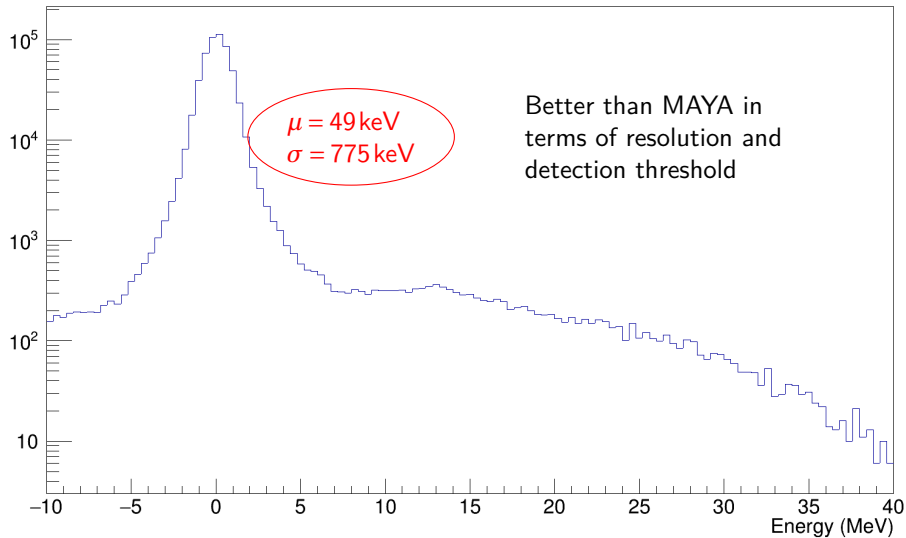
$^{58}\text{Ni}(\alpha, \alpha)^{58}\text{Ni}^*$ dataset

Kinematics: lab angle vs alpha energy



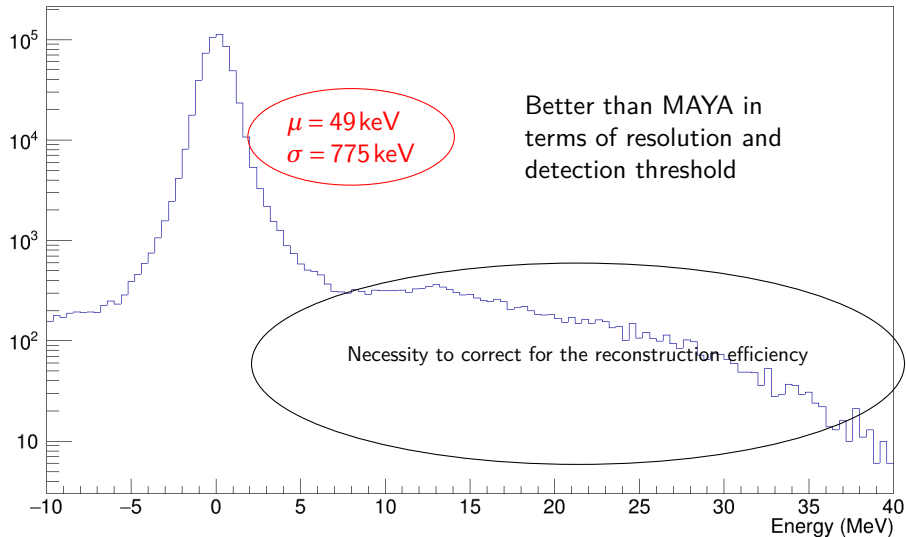
$^{58}\text{Ni}(\alpha, \alpha)^{58}\text{Ni}^*$ dataset

Kinematics: lab angle vs alpha energy

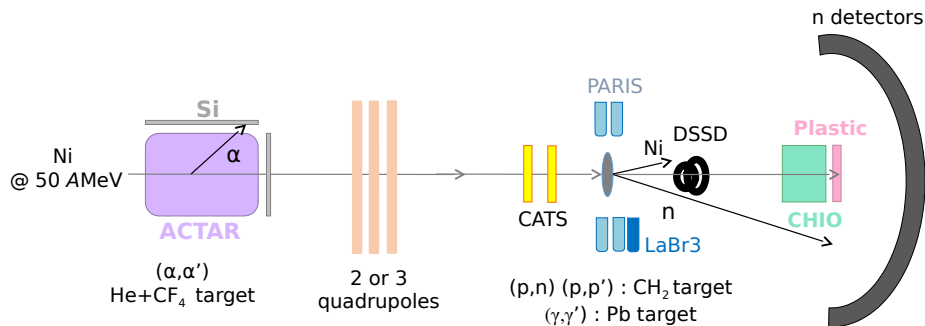


$^{58}\text{Ni}(\alpha, \alpha)^{58}\text{Ni}^*$ dataset

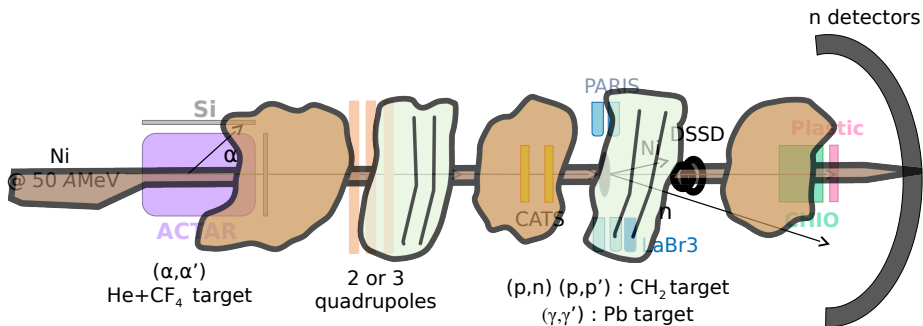
Kinematics: lab angle vs alpha energy



“Yakitori” (or “Brochette” [“Skewer”]) mode



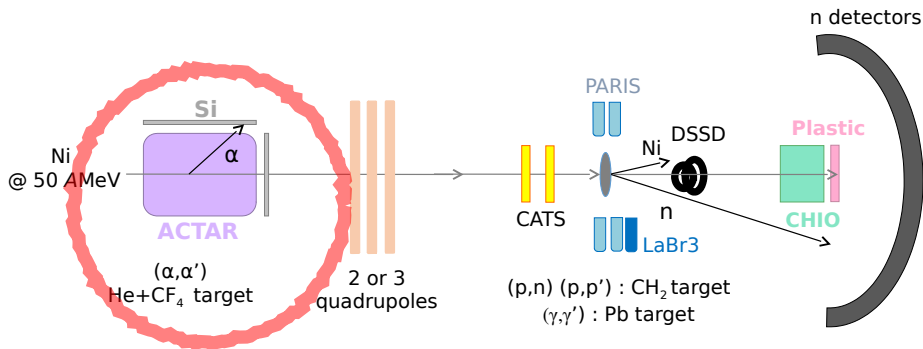
“Yakitori” (or “Brochette” [“Skewer”]) mode



Letter of Intent/Proposal

“Yakitori” (or “Brochette” [“Skewer”]) mode

GMR + ISGDR w/ ACTAR

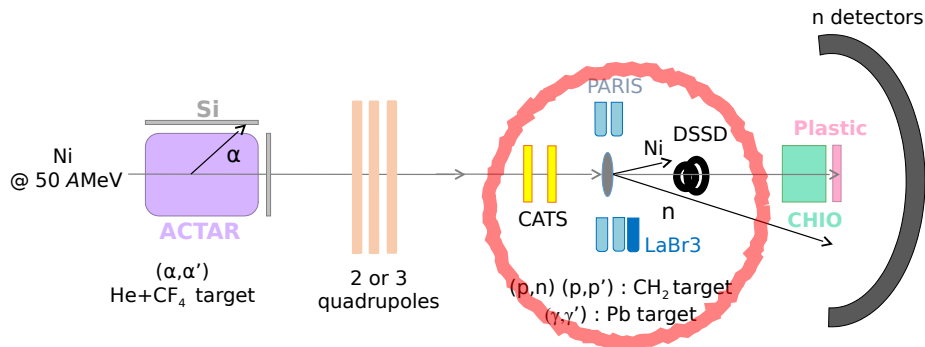


Letter of Intent/Proposal

“Yakitori” (or “Brochette” [“Skewer”]) mode

GMR + ISGDR w/ ACTAR

ISGDR + IVGDR w/ solid target & γ -rays

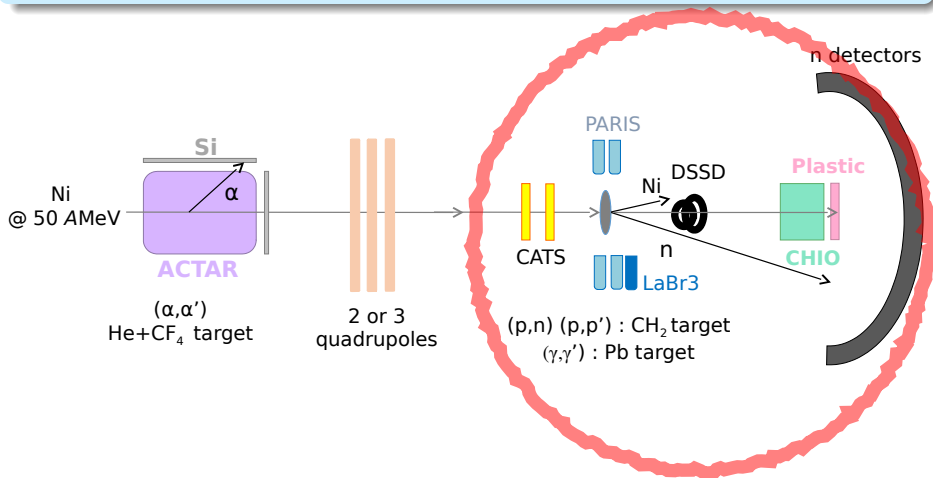


Letter of Intent/Proposal

“Yakitori” (or “Brochette” [“Skewer”]) mode

GMR + ISGDR w/ ACTAR

ISGDR + IVGDR w/ solid target & γ -rays



(IV) Pygmy Dipole Resonance

Illustration

GDR

PDR

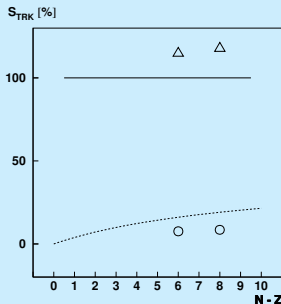
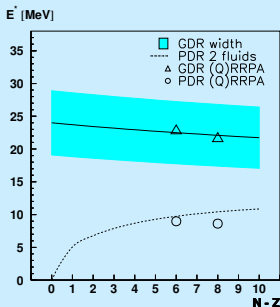
(IV) Pygmy Dipole Resonance

Illustration

GDR

PDR

Neon isotopes

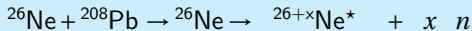
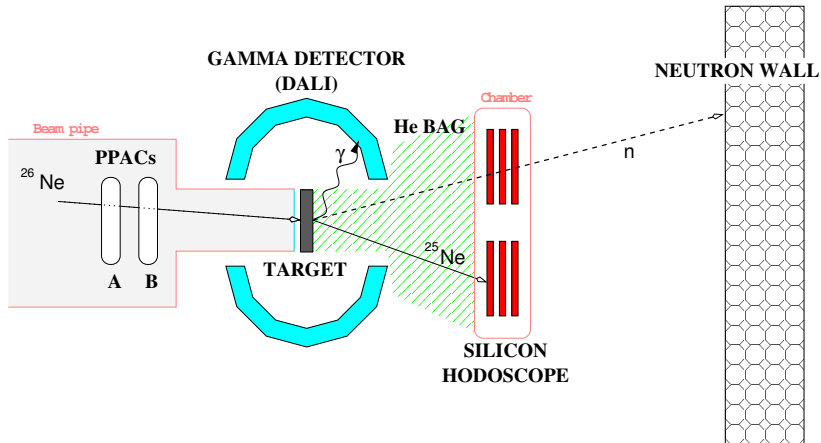


Suzuki et al. *Progress in Theoretical Physics* 83 (1990)

Cao et al. *Physical Review C* 71 (2005)

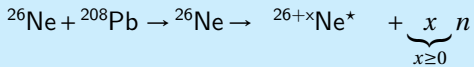
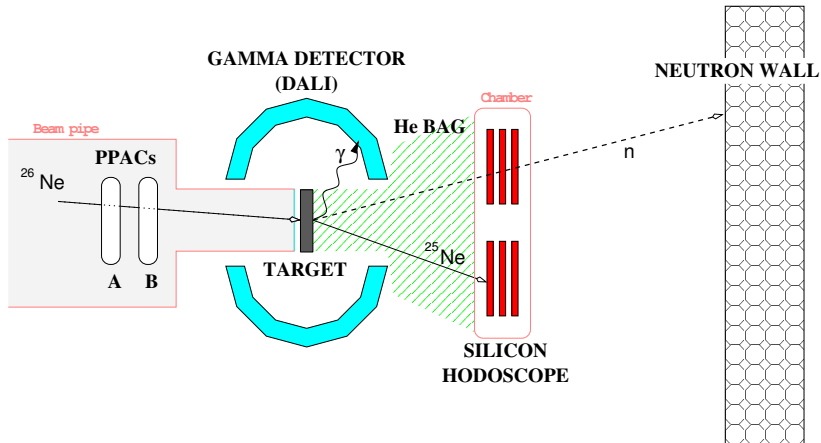
(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case



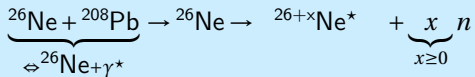
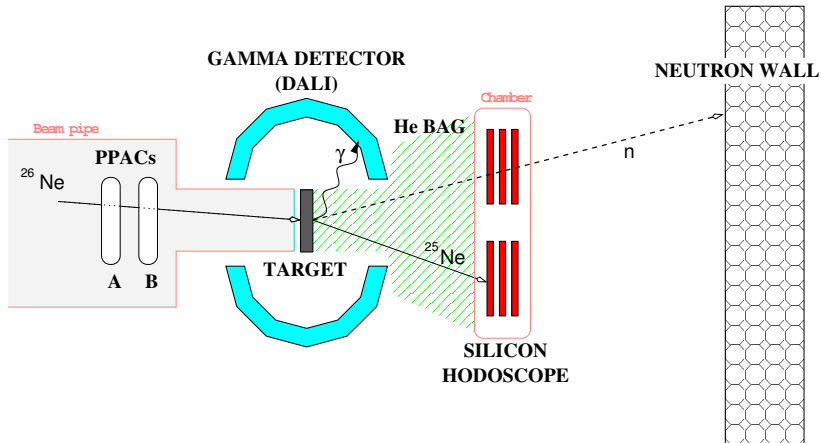
(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case



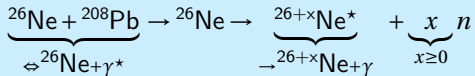
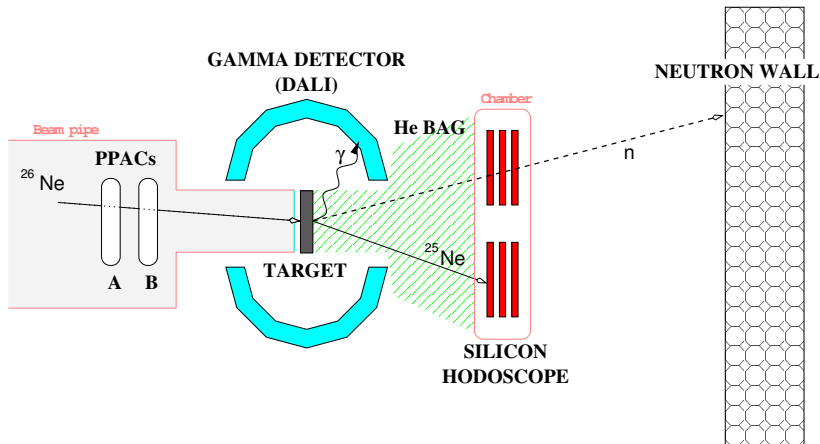
(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case



(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case



(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case

Invariant mass method:

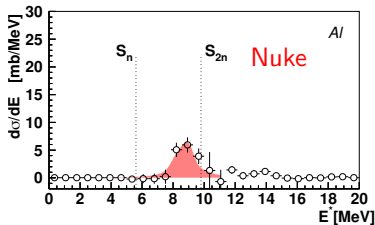
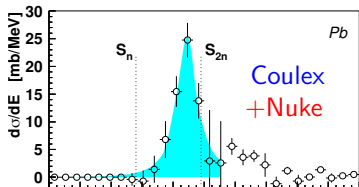
$$M_{\text{inv}} = M_0 + E^* = \sqrt{\sum_i E_i^2 - \sum_i \vec{p}_i^2} \Rightarrow E^* \sim E_{\text{rel}}(^{25}\text{Ne}, n) + E_{\gamma}(^{25}\text{Ne})$$

(IV) Pygmy Dipole Resonance

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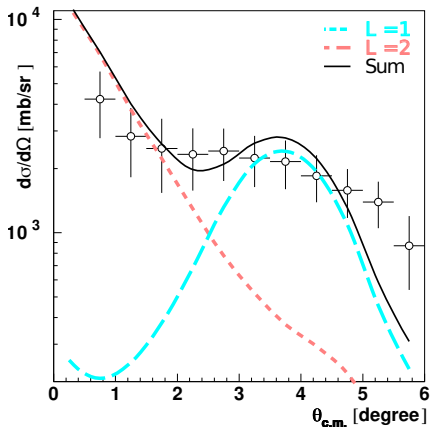
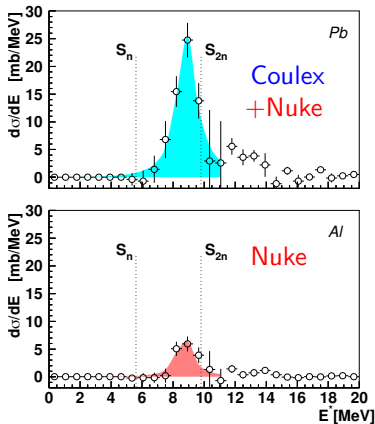


(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case

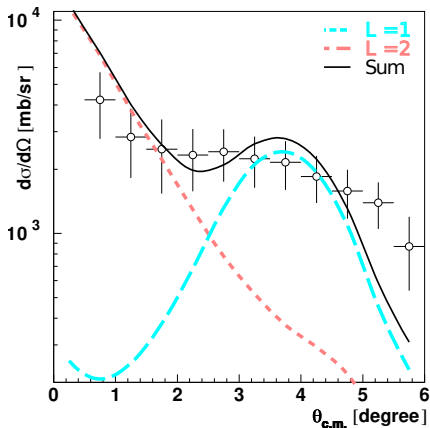
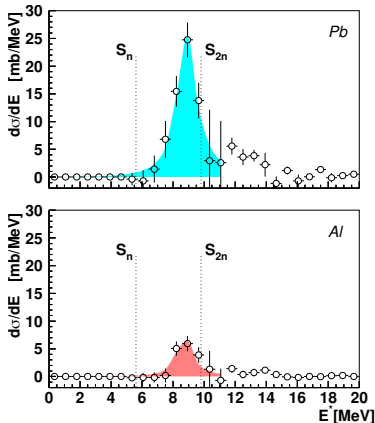
Invariant mass method:

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(IV) Pygmy Dipole Resonance

$^{26}_{10}\text{Ne}$ case



Results

$E_x \sim 9\text{MeV}$ and $S_{\text{TRK}} \sim 5\%$

Pygmy properties vs skin thickness

Results

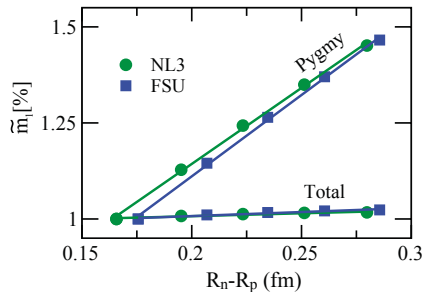
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Pygmy properties vs skin thickness

Results

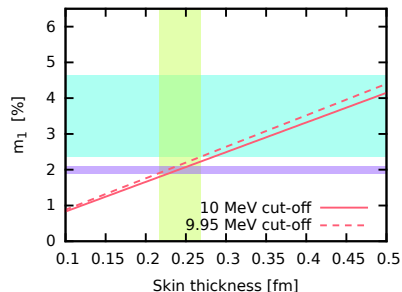
$E_x \sim 9\text{MeV}$ and $S_{\text{TRK}} \sim 5\%$

^{68}Ni



Piekarewicz *Phys. Rev. C* 83 (2011)

^{26}Ne



Inakura et al. *Phys. Rev. C* 88 (2013)

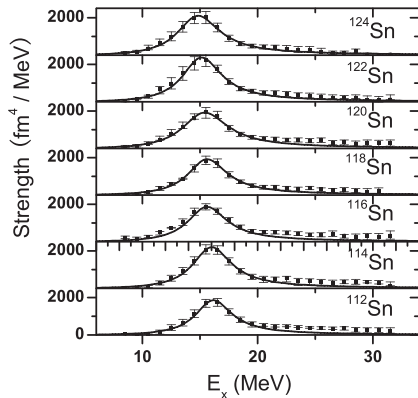
Marinova et al. *Phys. Rev. C* 84 (2011)

Gibelin PhD (2018)

Other setup for GMR measurement

Systematics in stable Sn isotopes

RCNP/Grand Raiden

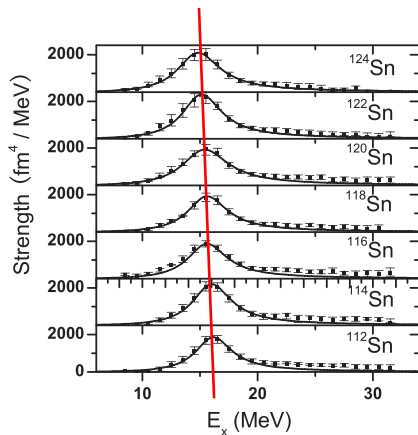


$$K_A - K_{\text{Coul}} Z^2 A^{-4/3} \sim K_{\text{vol}} (1 - A^{-1/3}) + K_{\tau} [(N - Z) / A]^2$$

$$K_{\text{Coul}} \sim \text{cst}$$

Systematics in stable Sn isotopes

RCNP/Grand Raiden

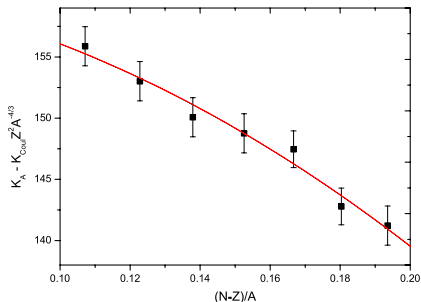
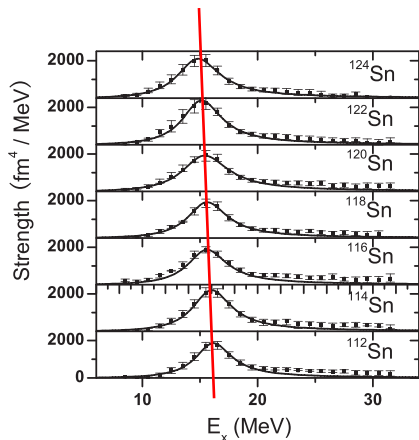


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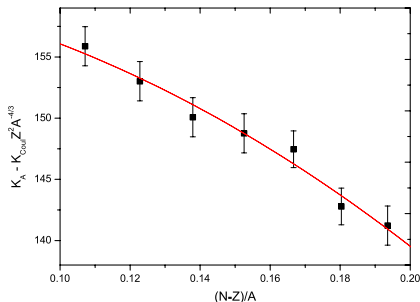
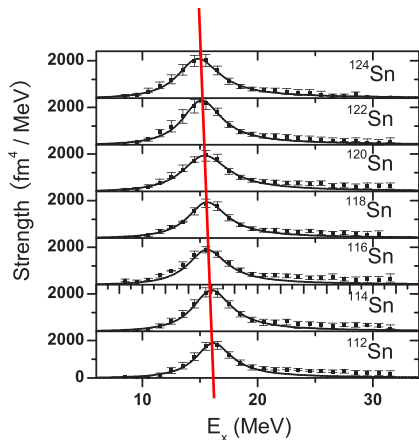


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Systematics in stable Sn isotopes

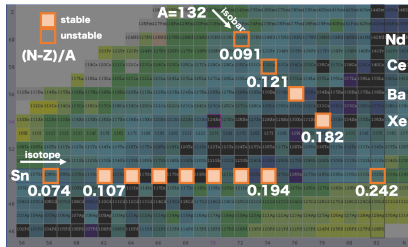
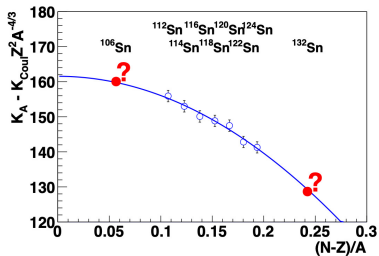
RCNP/Grand Raiden



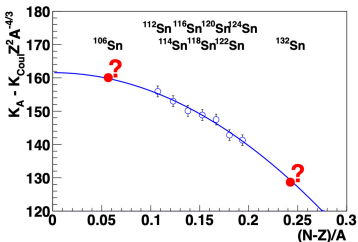
$$K_A - K_{\text{Coul}} Z^2 A^{4/3} \sim K_{\text{vol}} (1 - A^{-1/3}) + K_{\tau} [(N-Z)/A]^2$$

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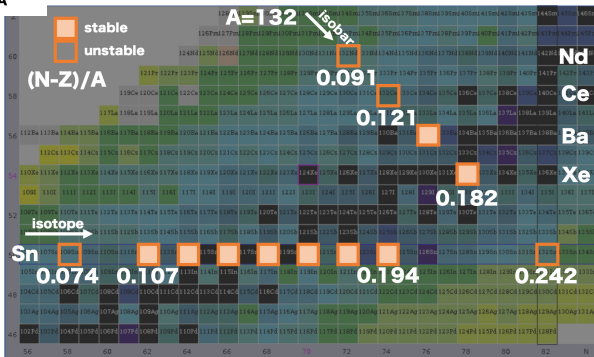
What about the radioactive Sn?



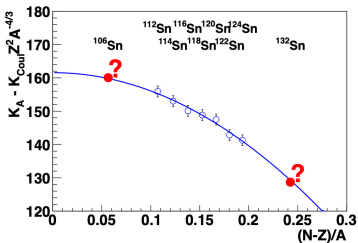
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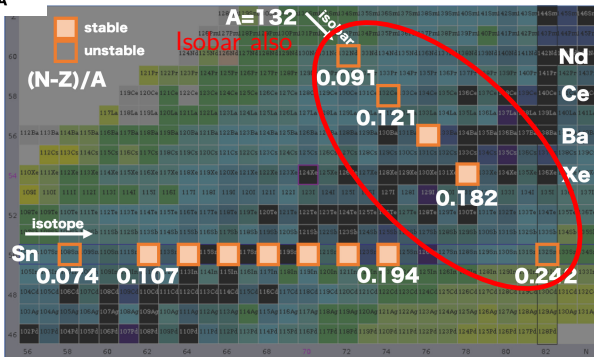
Combined measurement (RIBF + RCNP)
w/ the same setup "CATS"



What about the radioactive Sn?

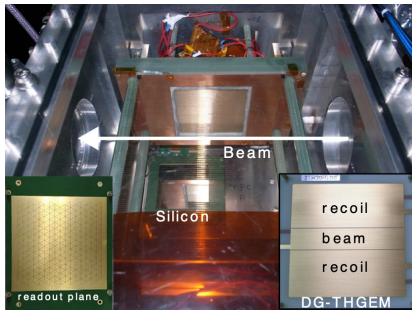
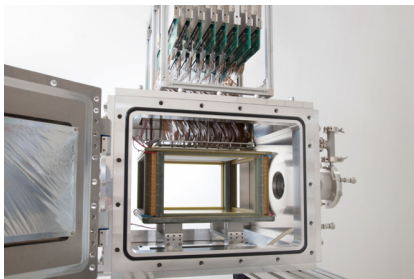


Combined measurement (RIBF + RCNP)
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CAT (CNS Active Target)

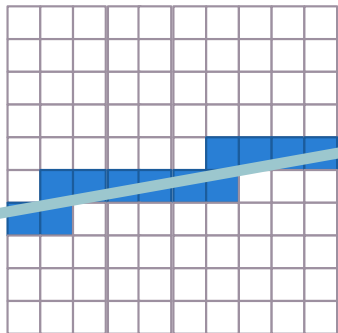
NEAT! (yEt another Active Target!)



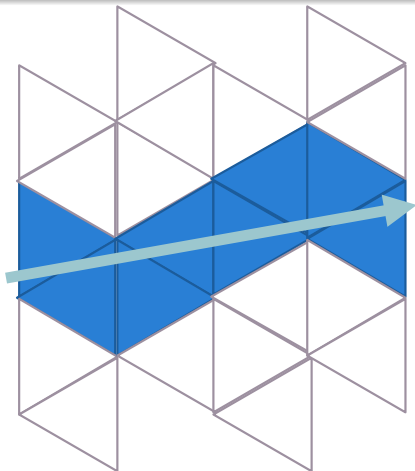
CAT (CNS Active Target)

NEAT! (yEt another Active Target!)

Concept: small number of readout channel (relatively low definition)



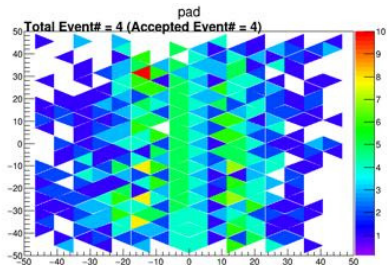
$$4 \text{ mm}^2 * 100 = 400 \text{ mm}^2 \text{ (4K)}$$



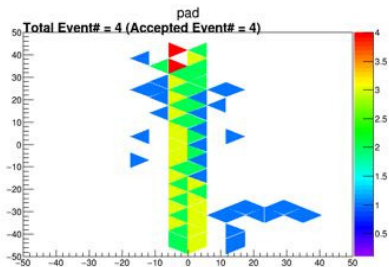
$$15.6 \text{ mm}^2 * 24 = 374 \text{ mm}^2 \text{ (HD) GEM technology}$$

CAT (CNS Active Target)

NEAT! (yEt another Active Target!)



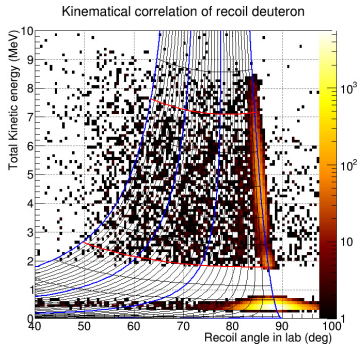
(RAW)



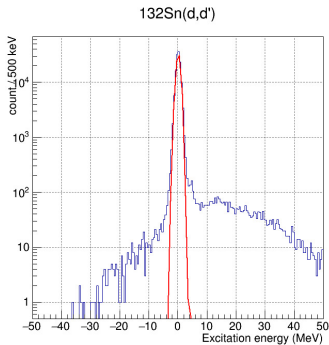
(CLEAN \leftarrow Hough transform (see Radon tr.))

CAT (CNS Active Target)

(preliminary) Results $^{132}\text{Sn}(d,d')$ (Courtesy of S. OTA [RCNP/Japan])



Kinematics

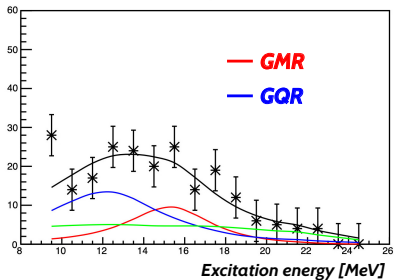


Excitation energy

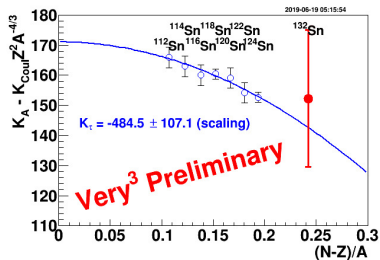
$$E_x \sim 0 \text{ MeV} (\sigma = 0.6 \text{ MeV})$$

CAT (CNS Active Target)

(preliminary) Results $^{132}\text{Sn}(d,d')$



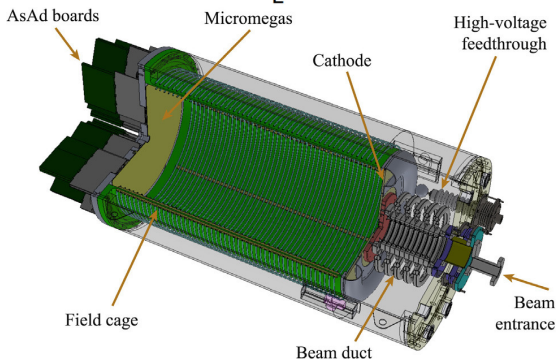
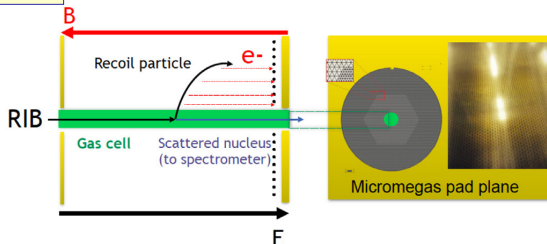
Lorentzian fit (GMR + GQR)



Compressibility

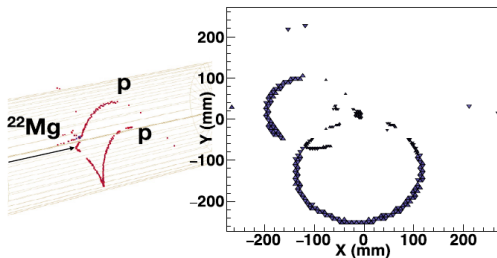
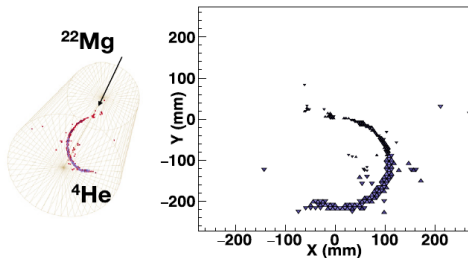
AT-TPC (Active Target Time Projection Chamber) – MSU

Ayyad et al. Nucl. Phys. A 954 (2020)



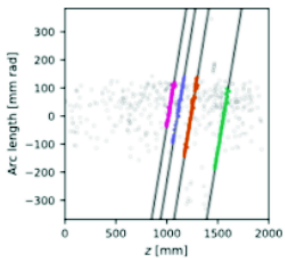
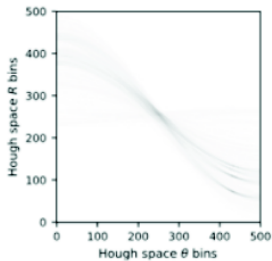
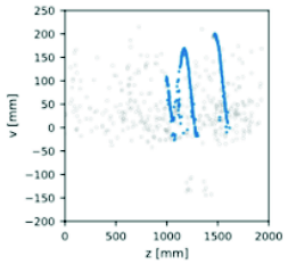
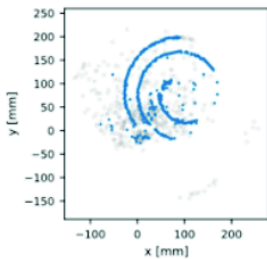
AT-TPC (Active Target Time Projection Chamber) – MSU

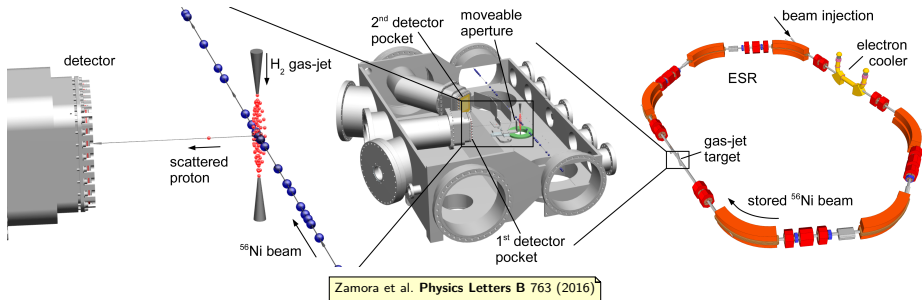
Ayyad et al. *Nucl. Phys. A* 954 (2020)



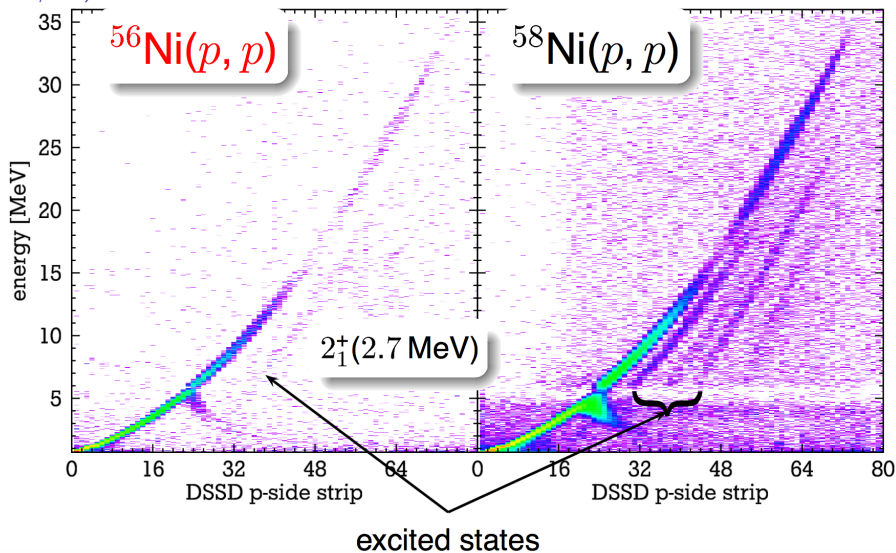
AT-TPC (Active Target Time Projection Chamber) – MSU

Ayyad et al. Nucl. Phys. A 954 (2020)





(FAIR/GSI)



Zamora et al. *Physics Letters B* 763 (2016)

- 1 Giant Resonances Properties
- 2 Experimental considerations
- 3 Example of experiments with stable nuclei
- 4 Example of experiments with unstable nuclei
 - GMR (MAYA/ACTAR)
 - PDR (Coulomb excitation)
- 5 Other setup for GMR measurement
 - RIKEN/RCNP
 - MSU
 - EXL (FAIR/GSI)

